



Nanofluid effect on forced convective heat transfer inside a heated horizontal tube

Mariem Elahmer^{1*}, Said Abboudi¹, Noureddine Boukadida²

¹ ICB UMR 6303, CNRS, Univ. Bourgogne Franche-Comté, UTBM, Département COMM, F-90010 Belfort, France

² Laboratory of Metrology and ENIM, 5000 Monastir, Tunisia

Email: mariem.elahmer@utbm.fr

ABSTRACT

In this paper, an unsteady forced convection of two-dimensional flow of Ethylene Glycol/CNT and Ethylene Glycol/CNT–Ag in horizontal cylinder is studied. The wall is submitted to a uniform or periodic heat flux. The transient distribution of the temperature is obtained by solving, at each axial position and time, a tridiagonal system by TDMA algorithm. The convective heat transfer coefficient and bulk temperature of nanofluid are investigated versus space-time and versus different nanoparticles volume fraction. The thermal phase shift is brought out for Ethylene Glycol base fluid and EG-CNT nanofluid. The obtained results show an improving of heat transfer rate when the volume fraction increases from 0 % to 10 % of CNT. Ethylene Glycol/CNT-AG seems to be more advantageous for the enhancement of convective heat transfer than EG/CNT. The results found can be exploited for various applications especially for the enhancement performances of heat transfer rate for the absorber tube of solar collector.

Keywords: Forced Convection, Laminar Flow, Unsteady, Hybrid Nanofluid, Conjugated Heat Transfer.

1. INTRODUCTION

Fluids such as water, Ethylene Glycol and oils have low thermophysical properties [1]. So, nanofluids have been exhibited novel and improved heat transfer due to their nanoscale size. For this reason, in the past years many experiments have been conducted numerically and analytically to validate the importance of nanofluids as working fluid. Almost all published papers in nanofluids field until nowadays studied the effect of concentration, shape, size, types of nanoparticles and its influence on heat transfer [2-3-4-5]. Maxwell [6-7] was the first who showed the possibility of enhancement of thermal conductivity of a solid-liquid mixture by increasing the volume fraction ratio of solid particles. In fact, Choi and Estman [8] presented for the first time a suspension of nanoparticles titled nanofluids and its benefit on different thermal systems to improve the heat transfer rate.

As a result, the first study on convective heat transfer of nanofluids in a circular tube was presented by Pak and Cho [9]. The results show that the Nusselt number increases with the increasing of volume fraction of nanoparticles. They proposed also a new correlation for the turbulent flow. Heris et al [10-11] investigated experimentally the convective heat transfer of CuO-water and Al₂O₃-water nanofluids with different concentrations in annular tube and laminar flow

under a constant wall temperature boundary condition. Also, Heris et al [12] studied numerically laminar convective heat transfer of nanofluid with constant wall temperature boundary condition through a circular tube. The results show that the heat transfer enhancement appears to be more pronounced with increase of nanoparticle volume fraction and the Nusselt number decreases with the increasing of nanoparticle size. The comparison of the numerical and experimental data shows good agreement.

Many researchers have focused on how to improve the efficiency of the solar collector experimentally. Onticar [13] has investigated experimentally the effect of different nanofluids as an absorbing medium on the efficiency of the micro solar thermal collector. The increased efficiency was 5%. Yousefi et al [14] has studied experimentally the effect of Al₂O₃-water nanofluid on the efficiency of a flat-plate solar collector. The results show that an efficiency improvement up to 28.3% for 0.2 wt% volume fraction of nanoparticles in nanofluid. Lenert and Wang [15] evaluated the influence of different variations in nanofluid volumetric receivers both experimentally and theoretically. Their results show that the efficiency of nanofluid volumetric receivers is increased with increasing solar concentration and nanofluid high.

Maiga et al [16] investigated numerically the thermal performance of nanofluid inside a uniformly heated tube. The flow is assumed as turbulent. The results show the heat transfer with the presence of nanofluid increase. A new correlation of Nusselt number is proposed to calculate the fully developed heat transfer coefficient for the nanofluid.

Maiga et al [17] studied numerically forced laminar convection flow using water/Al₂O₃ and Ethylene Glycol/Al₂O₃ nanofluids. Results show that the Ethylene Glycol/Al₂O₃ nanofluid has a higher heat transfer enhancement than water/Al₂O₃. So, the Ethylene Glycol is more advantageous than water. M.Izadi et al [18] studied numerically convective heat transfer of water/Al₂O₃ nanofluid in horizontal annulus where the wall is submitted to a uniform heat flux. Their results demonstrate that convection heat transfer coefficient and the nanofluid bulk temperature significantly increase with increasing of the nanoparticle volume fractions. Using the two-phase mixture model, Labib et al [19] developed a forced convection heat transfer study in which they showed that Ethylene Glycol/CNT-Al₂O₃ which gives better heat transfer enhancement than water/CNT-Al₂O₃. Moghadassi et Al [20] studied numerically the effect of Water/ Al₂O₃ nanofluid and Water/ Al₂O₃-Cu hybrid nanofluid on laminar forced convection in a horizontal circular tube. The volume concentration and the average particle size are fixed at 0.1% and 15nm respectively. Results show that the highest convective heat transfer coefficient is obtained for hybrid nanofluid. It was found that:

$$\left(\overline{Nu}\right)_{water/Al_2O_3-Cu} > \left(\overline{Nu}\right)_{water/Al_2O_3} > \left(\overline{Nu}\right)_{water}$$

Most of precedent studies were brought out in steady regime. Among studies that investigate transient forced convective heat transfer in tube, we cite the work of M.Fakoor et al [21]. They performed a new analytical model to predict laminar forced convection heat transfer inside a tube under arbitrary cyclic time dependent heat flux. One of their objectives is to determine the transient fluid flow response under a dynamically varying heat flux. Their results are verified successfully with obtaining numerical results.

Through this literature review, it is noted that studies of unsteady conjugate forced convection flow with nanofluids are scarce in literature. In our work, we focus on numerically unsteady laminar forced convection flow of nanofluids in a thick circular tube with an imposed parietal heat flux.

The main purpose of this study is to probe the effects of nanofluid: EG-CNT and hybrid nanofluid EG/CNT-Ag on conjugated convective heat transfer inside a heated horizontal thickness tube with uniform and harmonic heat flux densities. The heat transfer coefficient, and temperature field for the different studied cases, are presented in this work.

2. MATHEMATICAL MODELING AND NUMERICAL METHOD

Figure 1 shows the geometrical configuration of the studied problem. It consists of a two-dimensional circular tube with thickness wall.

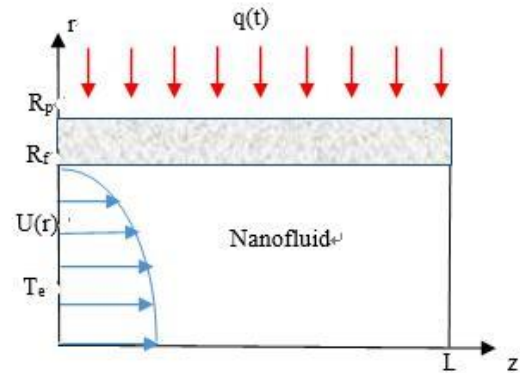


Figure 1. Schematic model

In this work, the nanofluid is assumed to be incompressible and Newtonian. Based on the literature [22], in the temperature and volume fraction ranges used, the nanofluids (EG-CNT) and (EG, CNT-Ag) are each considered as homogeneous and isotropic media, which makes it possible to assume, among other things, a single phase. The flow is unsteady, axisymmetric, laminar and fully developed. The thermophysical properties of the nanofluid are considered constant. The base fluid and the nanoparticle are in thermal equilibrium and have the same velocity. Effects of radiation and the viscous dissipation are neglected. The inlet velocity profile is parabolic and the inlet temperature is uniform.

Based on the previous assumptions, the transient conjugated heat transfer in the fluid and the solid tube is governed by the following equations and boundary conditions:

In the fluid $0 < r < R_f$

$$\rho_f C p_f \left(\frac{\partial T_f}{\partial t} + 2U_{max} \left(1 - \frac{r^2}{R_f^2}\right) \frac{\partial T_f}{\partial z} \right) = \lambda_f \left[\frac{1}{r} \frac{\partial T_f}{\partial r} + \frac{\partial^2 T_f}{\partial r^2} + \frac{\partial^2 T_f}{\partial z^2} \right] \quad (1)$$

In the solid $R_f < r < R_p$

$$\rho_s C p_s \left(\frac{\partial T_s}{\partial t} \right) = \lambda_s \left[\frac{1}{r} \frac{\partial T_s}{\partial r} + \frac{\partial^2 T_s}{\partial r^2} + \frac{\partial^2 T_s}{\partial z^2} \right] \quad (2)$$

2.1 Boundary conditions

For the fluid: $0 < r < R_f$

$$T_f(r, 0, t) = T_e, \quad \frac{\partial T_f}{\partial z}(r, L, t) = cst \quad (3a)$$

For the solid: $R_f \leq r < R_p$

$$\frac{\partial T_s}{\partial z}(r, 0, t) = 0, \quad \frac{\partial T_s}{\partial z}(r, L, t) = cst \quad (3b)$$

$$-\lambda_s \frac{\partial T_s}{\partial r}(R_s, z, t) = q(z, t) + h_e (T_a - T_s(R_s, z, t)) \quad (3c)$$

$0 < z < L$

$$q(z, t) = q_0 \text{ or } q(z, t) = q_0(1 + \sin(\alpha t))$$

$$\text{Interface solid-fluid: } T_s(R_f, z, t) = T_f(R_f, z, t),$$

$$\lambda_s \frac{\partial T_s}{\partial r}(R_f, z, t) = \lambda_f \frac{\partial T_f}{\partial r}(R_f, z, t), 0 < z < L \quad (3d)$$

$$\text{Initial conditions: } t=0, 0 \leq r < R_p, 0 < z < L$$

$$T_s(r, z, 0) = T_f(r, z, 0) = T_0 \quad (3e)$$

2.2 Thermophysical properties of nanofluid and hybrid nanofluid

The thermophysical properties of nanofluid are defined as follows [23]:

- Density and heat capacity

$$G_{nf} = (1 - \phi)G_f + \phi G_{np}, G = \rho, C_p \quad (4)$$

- Thermal conductivity of Hamilton & Crosser

$$\lambda_{nf} = \lambda_f \frac{\lambda_{np} + (n-1)\lambda_f + (n-1)(\lambda_{np} - \lambda_f)\phi}{\lambda_{np} + (n-1)\lambda_f - (\lambda_{np} - \lambda_f)\phi} \quad (5)$$

where the number n is an empirical shape factor: $n=3$ for spherical particles and $n=6$ for cylindrical particles.

- Dynamic viscosity of Einstein

$$\mu_{nf} = \mu_f(1 + 2.5\phi) \quad (6)$$

- Dynamic viscosity of Brinkman

$$\mu_{nf} = \frac{\mu_f}{(1 - \phi)^{2.5}} \quad (7)$$

- Dynamic viscosity of Batchelor

$$\mu_{nf} = \mu_f(1 + 2.5\phi + 6\phi^2) \quad (8)$$

The thermophysical properties of hybrid nanofluid are defined as follows [24]:

- Density and heat capacity: $G = \rho, C_p$

$$G_{hnf} = (1 - \phi)G_f + \phi_{np1}G_{np1} + \phi_{np2}G_{np2} \quad (9)$$

Thermal conductivity of Maxwell

$$\lambda_{hnf} = \lambda_f \frac{\frac{\phi_{12}}{\phi} + 2\lambda_f - 2\phi\lambda_f + 2\phi_{12}}{\frac{\phi_{12}}{\phi} + 2\lambda_f + \phi\lambda_f - \phi_{12}} \quad (10)$$

$$\phi_{12} = \phi_{np1}\lambda_{np1} + \phi_{np2}\lambda_{np2}$$

- Dynamic viscosity of Einstein

$$\mu_{hnf} = \mu_f(1 + 2.5\phi) \quad (11)$$

$$\phi = \phi_{CNT} + \phi_{Ag}, \phi_{CNT} = 9\% \text{ and } \phi_{Ag} = 1\%$$

2.3 Numerical procedure

The Fourier series method seems less adapted to problem of conjugate transfer and non-linear system. As a result, the finite difference method was considered to solve the system of equation (1), (2) and (3). For the purpose of numerical stability, a fully implicit formulation is adopted. The unsteady of energy storage and advection terms are represented by backward and upwind differences respectively and the central difference form is used to represent the axial and the radial diffusion terms. The transient distribution of the temperature is obtained by solving, at each axial position and time, a tridiagonal system by TDMA algorithm.

The choice of time step and space grid following the axial direction must satisfy the stability criterion of Von Neumann:

$$\frac{U_{\max} \Delta t}{\Delta z} < 1$$

2.4 Validation of the present code

Based on the strategy of developing our own codes in our research laboratory, a first FORTRAN code was developed. To test and verify it, we have adopted the same conditions used in the work of Ozan Sert et al. [25].

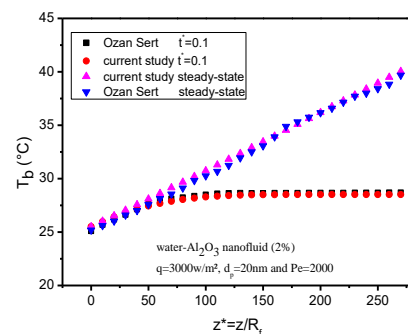


Figure 2. Comparison of the obtained transient and study state bulk temperatures with those of Ozan Sert et al. [25]

Figure 2 shows a comparison of the axial distribution of the bulk nanofluid temperature for the above adopted boundaries and initial conditions, with negligible wall thickness. Excellent agreement between the results is observed at transient and steady state regimes. Similarly, in the same conditions, Figure 3 shows a comparison of the convective heat transfer coefficient of nanofluid along the tube, without wall thickness. There is a small relative

difference between the current study and the values of Ozan Sert et al. [25] caused by their thermophysical properties which depend on temperature.

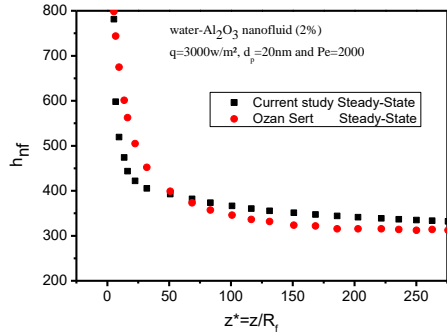
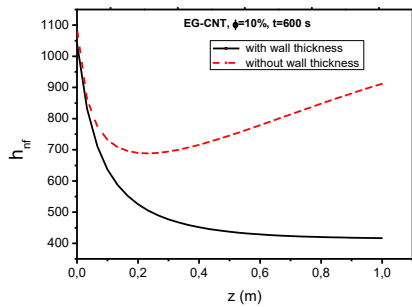
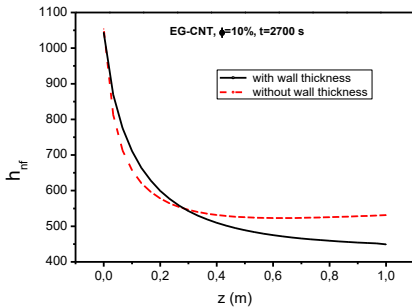


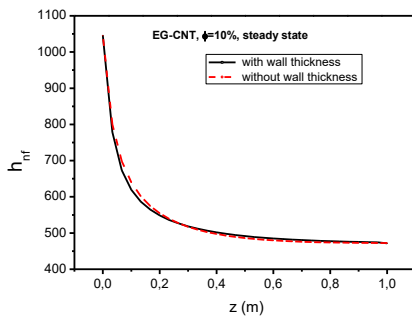
Figure 3. Comparison between the convective heat transfer coefficient of Ozan Sert et al. [25] and the current study



(a)



(b)



(c)

Figure 4. Evolution of the convective heat transfer coefficient: (a) at 600 s, (b) at 2700 s, (c) Steady State

Figure 4 shows the convective heat transfer coefficient along the tube with or without wall thickness at different times (t=600s: Figure 4a, t=2700s: Figure 4b, Steady state: Figure 4c) at a given nanoparticle volume fraction ($\phi = 10\%$).

As a result, the time has a significant effect on the convective heat transfer. At t=600s, fig 4a there is an important difference between the two cases: with and without wall thickness during the first instants of the transient regime, but at t=2700 s, fig 4b, it decreases as it approaches the steady state, fig 4c. This difference is principally due to the temperature gradient in the interface (nanofluid-steel wall) and the gap between the interfacial wall-fluid temperature and average bulk nanofluid one.

3. RESULTS AND DISCUSSION

In this study, unsteady laminar forced convection heat transfer of nanofluids in cylinder is investigated. Numerical simulations have been performed by using the following data:

- Geometry of the duct:
 $U_{\text{moy}} = 2.47 \cdot 10^{-2} \text{ m/s}, R_f = 5 \text{ mm}, R_s = 2 \text{ mm}$
- Physical properties of the duct:
 $\lambda_s = 40 \text{ W/mC}, \alpha_s = 3 \cdot 10^{-7} \text{ m/s}$
- Initial and boundary conditions:
 $T_0 = 293 \text{ K}, T_e = 293 \text{ K}, T_a = 293 \text{ K}, q_0 = 5000 \text{ W/m}^2,$
 $\omega = 3 \cdot 10^{-3} \text{ rd/s}, \text{Pe} = 1000.$
- -Discretization:
 $\Delta t = 1 \text{ s}, N_t = 7200, N_{r_f} = 51, N_{r_p} = 11, N_z = 31.$
- Physical properties of the nanofluid and hybrid nanofluid: see tables 1 and 2.

Table 1. Thermophysical properties of base fluid and nanofluid

	EG	10% EG/CNT
Kg.m^3	998,3	1038.5
$\text{W.m}^{-1}\text{K}^{-1}$	0.602	0.8026
$\text{C}_p \text{ J.kg}^{-1}\text{K}^{-1}$	4182.2	3901.98
E Pa.s^{-1}	$1,001 \cdot 10^{-3}$	$1,25 \cdot 10^{-3}$
Pr	6.954	4.868

Table 2. Thermophysical properties of nanoparticles and hybrid nanofluid

	CNT	Ag	EG/CNT-Ag
Kg.m^3	1400	10500	1133.5
$\text{W.m}^{-1}\text{K}^{-1}$	3500	429	0.825
$\text{C}_p \text{ J.kg}^{-1}\text{K}^{-1}$	1380	235	3862.5
Pa.s	-	-	$1.25 \cdot 10^{-3}$

3.1 Physical aspects

The volume fraction has an effect on many additional parameters like, thermal conductivity, dynamic viscosity and pressure drop. In this work, the effect of the pressure drop is

low because we supposed the flow is established. Figure 5 shows the variation of the Prandtl number with the nanoparticle volume fraction for two studied nanofluids (EG/CNT and EG/Al₂O₃). The choice of EG as base fluid is referred to the fact that it leads to a higher exchange of heat transfer as it was mentioned in literature. It is useful to recall that the EG/CNT Prandtl number is higher than EG/Al₂O₃ one, mostly for $\phi > 4\%$.

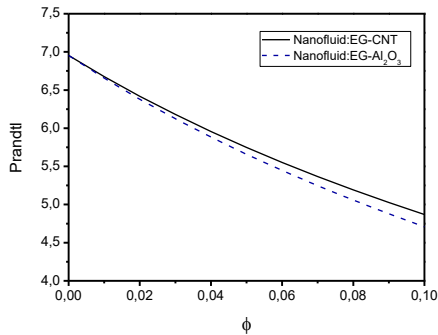


Figure 5. Variation of the Prandtl number with the nanoparticle volume fraction for two studied nanofluids (EG/CNT and EG/Al₂O₃)

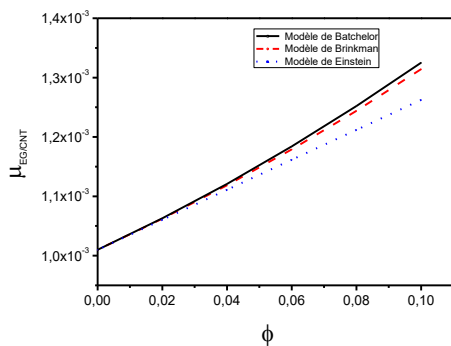


Figure 6. Dynamic viscosity of nanofluid for three models versus nanoparticle volume fractions

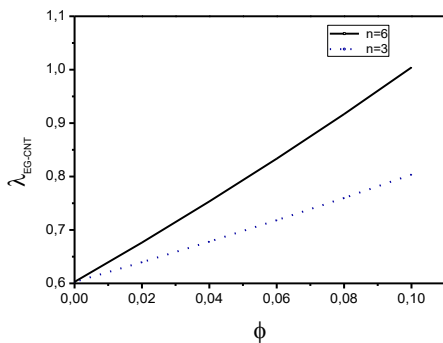


Figure 7. Thermal conductivity of nanofluid versus nanoparticle volume fractions for cylindrical ($n=6$) and spherical nanoparticles ($n=3$)

For Einstein, Brinkman and Batchelor classical models, Figure 6 shows the evolution of EG/CNT nanofluid dynamic viscosity versus nanoparticle volume fraction. As it is depicted, between $\phi = 0\%$ and $\phi = 4\%$, the increase of the dynamic viscosity is identical for all models. But for

$4\% \leq \phi \leq 10\%$, Einstein model which is retained in our study, gives lower values of μ compared to the two others models.

Figure 7 shows the effect of CNT nanoparticles shape on EG/CNT nanofluid thermal conductivity. It can be concluded that the shape of nanoparticles has an effect on the thermal conductivity. The cylindrical nanoparticles have higher conductivity than the spherical nanoparticles.

3.2 Thermal aspects

3.2.1 Uniform heat flux

3.2.1.1 Temperature profiles

The radial profile of the nanofluid temperature of EG/CNT at three axial positions ($z = L/3, 2L/3, L$) is reported at different times in Figures.8a and 8b.

The thickness of the wall and its thermal conductivity make that there is no significant effect on the radial variation on the temperature of the wall.

Figure 9 shows the development of temperature profiles of nanofluid and steel wall of the tube. Given the fact that the nanofluid receives continuously a heat flux by the tube wall, its temperature continues to increase with z . The nanofluid bulk temperature is much lower at short time than at the steady state regime. By following the evolution of the tube wall temperature as a function of z , it is concluded that $\frac{\partial T_s}{\partial z} \neq 0$ whatever the time. As it is depicted in Figure 10,

for a given time, the nanofluid bulk temperature increases with z . For a given location, it increases for $0 < t < t_e$ where t_e is the time at which the thermal regime becomes established. This time also increases with axial positions.

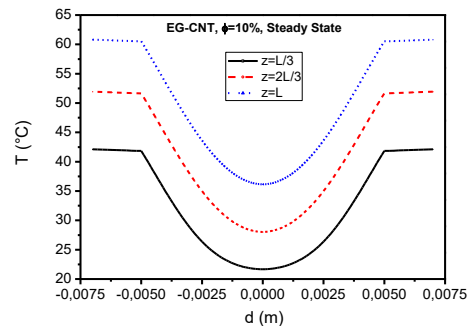


Figure 8a. Steady state radial profiles of temperature at different axial positions

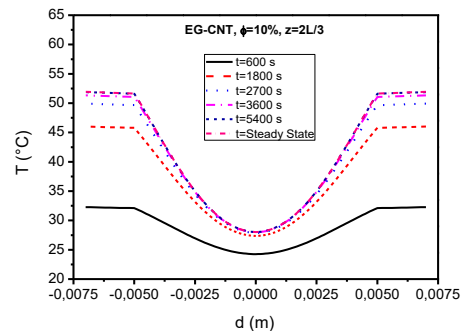


Figure 8b. Radial profiles of temperature at different instants

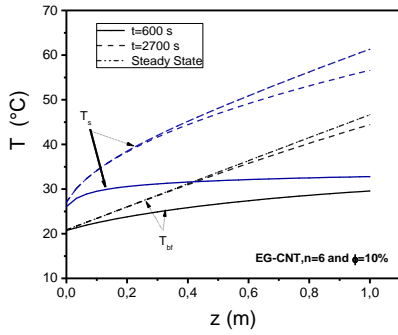


Figure 9. Transient distributions of bulk temperature of nanofluid and wall temperature

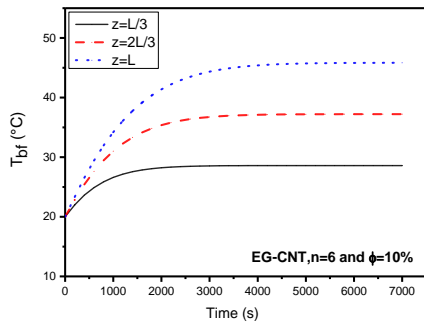


Figure 10. Evolution of bulk nanofluid temperature for different axial positions

3.2.1.2 Convective heat transfer coefficient profiles

Figure 11 shows the axial profile of the convective heat transfer coefficient for different times. It can be seen that the heat transfer coefficient slightly increases with time and decreases along the tube length. Also, as expected, a significant increase of the transfer coefficient is observed for EG/CNT nanofluid in comparison with that obtained with EG pure. This increasing is explained by the augmentation of the thermal conductivity.

Figure 12 illustrates a comparison between the effects of cylindrical particles and the spherical particles on convective heat transfer coefficient. Mostly for $z > 0.18$ m, it is clearly observed, that the cylindrical shape ($n=6$) leads to higher values of h_{nf} than those of spherical shape ($n=3$). This type of development is intimately linked to the thermal conductivity evolution shown in Figure 7.

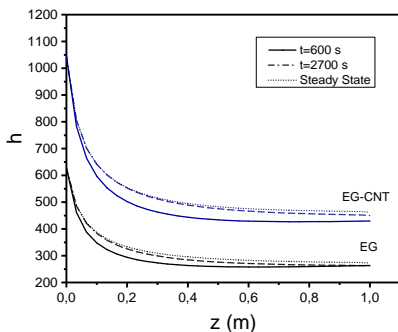


Figure 11. Transient axial profiles of the convective heat transfer coefficient for Ethylene Glycol and EG-CNT

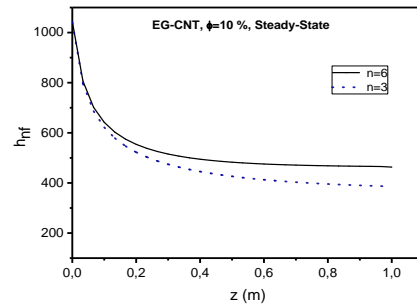


Figure 12. Convective heat transfer coefficient for cylindrical ($n=6$) and spherical nanoparticles ($n=3$)

For EG/CNT nanofluid, we present in Figure 13 the axial distribution of the convective heat transfer coefficient, for different nanoparticle volume fractions and in Figure 14, its evolution for the mentioned axial positions. As it is seen, this coefficient increases with ϕ , and the steady state regime is reached at time $t=5400$ s. From Figure 10, it is possible to conclude that nanofluid with higher nanoparticle volume fractions has more ability to exchange heat transfer. The values of the convective heat transfer coefficient slightly decrease with an increase in the axial position.

For $\phi=10\%$, Figure 15 presents a comparison between the convective heat transfer coefficient for EG/CNT and EG/CNT-Ag. As is expected, higher heat transfer coefficient is observed for hybrid nanofluid. This is due to the higher value of thermal conductivity of hybrid nanofluid compared to sample nanofluid.

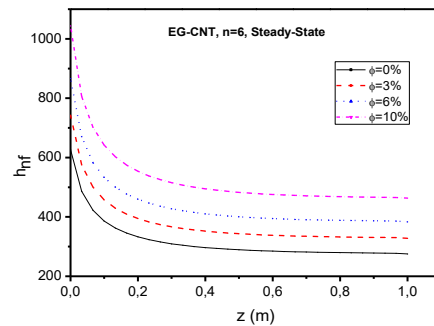


Figure 13. Convective heat transfer coefficient for different nanoparticle volume fractions

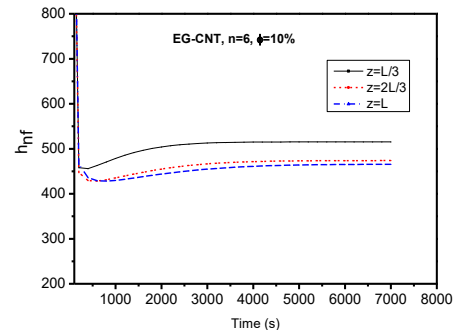


Figure 14. Evolution of heat transfer coefficient of EG-CNT nanofluid for different axial positions

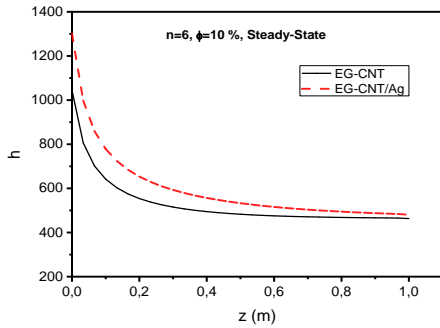


Figure 15. Distribution of convective heat transfer coefficient for EG-CNT and EG/CNT-Ag

3.2.2 Variable heat flux

The imposed wall density of heat flux is given by the following periodic form:

$$q(t) = q_0(1 + \sin(\omega t)), \omega = 3 \cdot 10^{-3} \text{ rad/s}, q_0 = 5000 \text{ W/m}$$

3.2.2.1 Temperature profiles

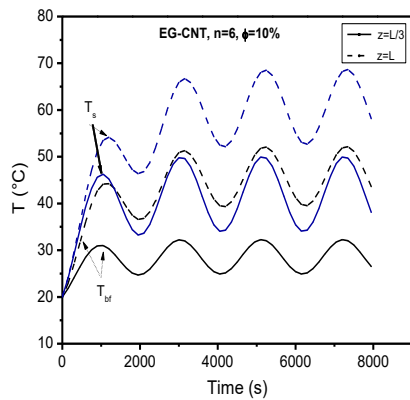


Figure 16. Variation of the temperature versus time at two axial positions

Figure 16 shows the evolution of the nanofluid bulk temperature and steel wall temperature versus time at two axial positions. It can be seen that the nanofluid bulk temperature oscillates with time and increases with z for a given time. At initial times, the values of the temperature slightly increase with axial position. After a time, response of one hour, the nanofluid bulk temperature oscillates with the quasi-same magnitude. It should be noted that the nanofluid bulk temperature shows a similar behavior as the tube wall temperature versus the time.

For different thicknesses and for $z=2L/3$, figure 17 shows the evolution of the average temperature of EG base fluid and EG-CNT nanofluid. We observe the presence of the reduction of thermal amplitude with the thickness, this is due to the role of the thermal inertia which has for main role to oppose the variation of the temperature.

It is concluded that whatever the thickness value the presence of nanoparticles in the base fluid causes a thermal phase shift.

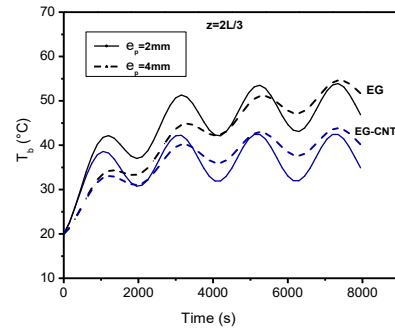


Figure 17. Evolution of the average temperature of EG and EG-CNT nanofluid for two wall thicknesses

3.2.2.2 Convective heat transfer coefficient profiles

Figure 18 shows the axial convective heat transfer coefficient for different volume fraction nanoparticles. The same conclusion, with an imposed constant flux, is observed. Figure 19 illustrates the evolution of the heat transfer coefficient of nanofluid with time at different axial positions. At the first instants, the values of the heat transfer coefficient are significantly higher and decrease rapidly. After that, they begin to oscillate with amplitudes which diminish gradually as one approaches the exit of the tube.

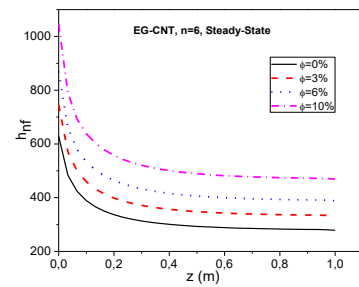


Figure 18. Convective heat transfer coefficient for different nanoparticle volume fractions

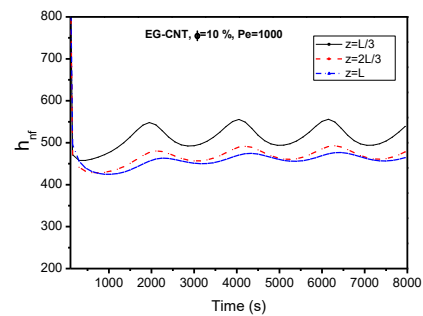


Figure 19. Evolution of heat transfer coefficient of EG-CNT nanofluid for different axial positions

3.2.2.3 Thermal phase shift profiles

The thermal phase shift between two temperatures peaks is expressed by $\Delta\varphi = \omega \Delta t$, where ω is the angular frequency of the imposed wall heat flux and Δt is the time difference between two peaks. It principally depends of the layer thickness and the thermophysical properties of the material.

Figure 20 shows the time evolution of interfacial temperature at $z=L$ for EG and EG-CNT nanofluid for two Peclet number,

Pe=1000 and Pe=2000 at $e_p=2$ mm, where the thermal phase shift is caused by the moving nanofluid and by steel layer of the tube.

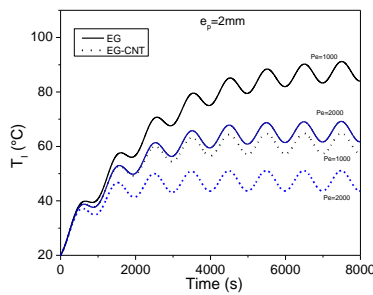


Figure 20. Time evolution of interfacial temperature at $z=L$ for EG and EG-CNT nanofluid

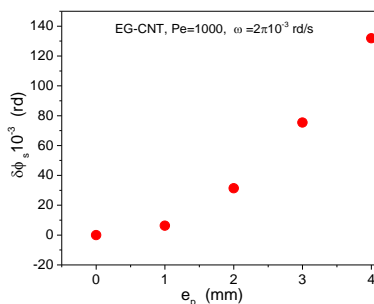


Figure 21. Wall thermal phase shift of EG-CNT nanofluid at different thickness

Figure 21 shows the wall thermal phase shift at different thickness of nanofluid at a fixed time $t=1000$ s and Peclet number $Pe=1000$. As it is depicted the thermal phase shift increased with the increasing of the wall thickness.

4. CONCLUSION

In this work, a numerical study of laminar conjugated heat transfer mechanisms of EG/CNT nanofluid and EG/CNT-Ag hybrid nanofluid inside a heated horizontal tube is proposed to improve the efficiency of an absorber tube of the solar collector. The main conclusions, from the discussed results, are drawn as follow:

- The dynamic viscosity for the three models (Einstein, Brinkman and Batchelor) and the thermal conductivity of Hamilton&Crosser model increases with the nanoparticles volume fractions.
- The shape of nanoparticles has an effect on the thermal conductivity which is higher for cylindrical particles than spherical particles.
- Whatever the wall heated flux nature (uniform or variable), convection heat transfer coefficient increases with nanoparticles volume fraction.
- Nanofluid bulk temperature increases with the axial direction and with time.
- For a given position, the heat transfer coefficient h_{nf} increases versus time.

-A phase shift appears in the presence of nanoparticles in the base fluid.

$$h_{EG/CNT-Ag} > h_{EG/CNT} > h_{pure EG}$$

This work has not been finished, so a continuous study on the radiation phenomenon of the solar collector seems to be necessary, to find a comparison between the efficiency of the solar collectors based water and based nanofluid.

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NOMENCLATURE

C _p	specific heat, J. kg ⁻¹ . K ⁻¹
d	diameter of the tube, m
d _p	diameter of the nanoparticle, m
e _p	e _p =R _p -R _f
h	heat transfer coefficient, W. m ⁻² .K ⁻¹ ,
U	velocity, m.s ⁻¹
L	length of the tube, m
n	shape factor
N	number of nodes
Nu	Nusselt number
Pe	Peclet number
q(t)	transient heat flux,
q ₀	constant heat flux, W.m ⁻²
r	radial coordinate, m
R _f	internal radius, m
R _p	external radius, m
T	temperature, K
T _a	ambient temperature, K
T _e	inlet temperature, K
T _s	wall temperature, K
T ₀	initial temperature, K
t	time, s
z	axial coordinate
EG	Ethylene Glycol
CNT	Carbon nanotube
Ag	Argent
Al ₂ O ₃	Alumina

Greek symbols

α	thermal diffusivity, m ² . s ⁻¹
φ	nanoparticle volume fraction
μ	dynamic viscosity, kg. m ⁻¹ .s ⁻¹
λ	thermal conductivity, W.K ⁻¹ .m ⁻¹
ρ	density, Kg.m ⁻³
ψ	nanoparticle sphericity
φ	thermal lag (phase shift)
ω	angular frequency

Subscripts

np	nanoparticle
f	fluid
s	solid
nf	nanofluid
e	established thermal regime