



MHD Flow and Heat Transfer of a Jeffrey Fluid over a Porous Stretching/Shrinking Sheet with a Convective Boundary Condition

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ABSTRACT

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This work explores the heat transfer flow characteristics of an incompressible non-Newtonian Jeffrey fluid over a stretching/shrinking surface with thermal radiation and heat source. The sheet is linearly stretched in the presence of a transverse magnetic field with convective boundary conditions. Appropriate similarity variables are used to transform the basic governing equations (PDEs) into ODEs. The resulting equations are solved by utilizing MATLAB *bvp4c*. The impact of distinctive physical parameters and dimensionless numbers on the flow field and heat transfer is analysed graphically. It is noticed that the measure of heat raised with increasing the Biot number and opposite effect with the rise of the suction parameter.

1. INTRODUCTION

In a real situation, many industrial fluids such as ketchup, pastes, slurries, paint, shampoo, blood, glues, printing inks, food materials, soap and detergent slurries, polymer solutions are non-Newtonian in nature. These fluids are basically non-linear and normally show both viscous and elastic properties. The constitutive equations relating such fluids are intrinsically more sophisticated than conventional Newtonian (Navier-Stokes) fluids. Most non-Newtonian models (Maxwell models, Oldroyd-B models, Walters-B short memory models, Jeffrey model and Eyring-Powell models etc.) involve varying degrees of refinement to the classical momentum conservation equations. Out of these, Jeffrey model is the simplest one rate type non-Newtonian liquids which exhibits shear diminishing attributes, yield pressure and high shear viscosity. Due to its simplicity and variety of applications in science and engineering, many researchers have been attracted. Kothandapani and Srinivas [1] contemplated the peristaltic flow of a Jeffrey liquid over an asymmetric channel in the occurrence of the magnetic field. Radiative effect on rheological fluid (Jeffrey fluid) over a stretchable cylinder in a 2D flow has been deliberated by Hayat et al. [2]. Khan et al. [3] presented the influence of third-grade nanofluid flow generated due to sheet stretching. Mishra et al. [4] implemented a shooting method to analyze the 2D flow of Jeffrey fluid through the stretchable sheet. Narayana et al. [5] discussed the viscous dissipation effects on Jeffrey fluid flow caused due to a stretchable sheet. More recently, many authors have studied the diverse non-Newtonian fluid flow models with various heat transfer effects [6-10].

The flow analysis of an electrically conducting non-Newtonian liquid over a shrinking/stretching surface plays a

vital role in mechanical and engineering applications such as metal spinning, hot rolling and shrinking film for packing of huge products etc. Shrinking surface flow is not quite same as that of stretching sheet flow due to the attraction of fluid towards the origin (slot). Hamid et al. [11] assumed a time-dependent magnetic field for analyzing the nature of Ohmic heating and viscous dissipation on Williamson fluid flow generated by stretching/shrinking sheet. Yasin et al. [12] considered the effect of thermal radiation on MHD flow past a shrinking/stretching sheet and solved the model by obtaining the dual solutions. Gangadhar et al. [13] investigated the effect of Newtonian heating on MHD micropolar nano-liquid in the case of stretching and shrinking surface. They dissected the existence of dual answers for the cases of stretching and shrinking. Jahan et al. [14] illustrated the influence of nanofluid over a stretching /shrinking sheet. Later, several authors discussed the effect of MHD flow on various fluid models [15-22].

The boundary layer flow with convective boundary conditions and heat transfer has much importance due to their applications in manufacturing and ecological technologies including energy storage, nuclear plants, gas turbines, geothermal reservoirs and rocket propulsion. In view of this, many investigators [23-28] concentrated the heat transfer and mass transfer concepts over various flow geometries. Zeeshan et al. [29] explored the Couette-Poiseuille flow of viscous nanofluid with convective boundary conditions by utilizing HAM technique. Shehzad et al. [30] analyzed the 3-D stretched flow of an incompressible Jeffrey liquid with convective conditions. Tlili et al. [31] applied thermodynamics theory to analyze the flow of nanofluid through a static wedge with convective boundary conditions. The impact of chemical reaction on MHD flow of Jeffrey nanofluid by utilizing the

convective boundary conditions was described by Kothandapani and Prakash [32]. The impact of the heat source on Jeffrey liquid over a stretching sheet in the presence of nonlinear energy and concentration has been focused by Narayana and Babu [33]. Recently, Ashraf et al. [34] and Mishra et al. [35] discussed the 2D flow of non-Newtonian convective flow with heat source/sink.

Motivated by these facts, a mathematical model on MHD boundary layer flow of a Jeffrey liquid with heat transfer owing to a permeable stretching sheet is proposed. Instead of frequently used uniform surface temperature, a convective boundary condition is applied which makes the analysis unique and further the consequences are accurate and practically valuable. An appropriate similarity transformation is applied to convert the basic governing equations (PDEs) into ODEs which are then solved by utilizing MATLAB bvp4c. The effects of various parameters on the flow field are shown graphically and discussed. It is believed that the present results not only find the applications in realistic engineering but also assist as a complement to the earlier research works.

2. MATHEMATICAL FORMULATION

A steady two-dimensional (x, y) boundary layer flow of an incompressible non-Newtonian Jeffrey liquid over a stretching sheet with thermal radiation and heat source is analysed. In the Cartesian coordinate system, the linear velocity of stretching sheet as $u_w(x) = bx$ (where b is a real number) is toward the x -axis and y -axis is opposite to it [see Figure 1]. Also, it is considered that the stretching surface is heated due to convection from a hot liquid at temperature T_f . Additionally, a uniform magnetic field of strength B_0 is forced normal to the plate along the y -axis and the induced magnetic field is neglected because the magnetic Reynolds number is very small. The expressions for the Cauchy and the extra stress tensors τ and S_1 in a Jeffrey non-Newtonian liquid are defined as follows [1]:

$$\tau = -pI + S_1 \quad (1)$$

$$S_1 = \frac{\mu}{1+\lambda} \left[A_1 + \lambda_1 \left(\frac{\partial}{\partial t} + V \cdot \nabla \right) A_1 \right] \quad (2)$$

In which p is the pressure. The Rivlin-Ericksen tensor A_1 is defined as $A_1 = \nabla V + (\nabla V)'$, where V is the velocity field vector.

The governing equations of non-Newtonian liquid under the assumption of boundary layer approximations is given by (Ref. [8, 23]).

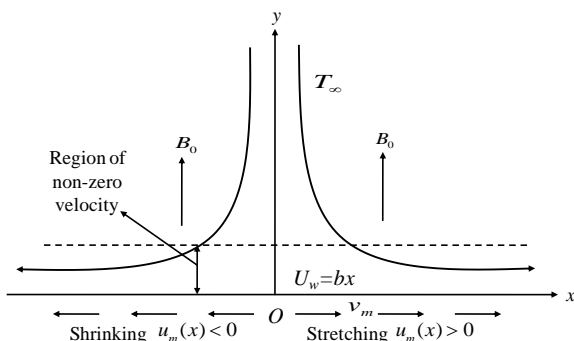


Figure 1. Geometry of the problem

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0 \quad (3)$$

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = \frac{\nu}{1+\lambda} \left\{ \frac{\partial^2 u}{\partial y^2} + \lambda_1 \left[\frac{\partial u}{\partial y} \frac{\partial^2 u}{\partial x \partial y} + u \frac{\partial^3 u}{\partial x \partial y^2} \right] - \frac{\partial u}{\partial x} \frac{\partial^2 u}{\partial y^2} + v \frac{\partial^3 u}{\partial y^3} \right\} - \frac{\nu}{K^*} u - \frac{\sigma B_0^2}{\rho} u \quad (4)$$

$$u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \alpha^* \frac{\partial^2 T}{\partial y^2} - \frac{1}{\rho c_p} \frac{\partial q_r}{\partial y} + \frac{Q_0}{\rho c_p} (T - T_\infty) \quad (5)$$

The boundary conditions:

$$\left. \begin{aligned} u &= u_w(x) = bx, v = v_w, \\ -k \frac{\partial T}{\partial y} &= h(T_f - T) \text{ at } y = 0 \\ u &= 0, u' \rightarrow 0, T \rightarrow T_\infty \\ &\text{as } y \rightarrow \infty \end{aligned} \right\} \quad (6)$$

Utilizing the Rosseland approximation, the radiative heat flux q_r is defined as [36]:

$$q_r = -\frac{4\sigma^*}{3K_s} \frac{\partial T^4}{\partial y} \quad (7)$$

where, K_s and σ^* is the mean absorption coefficient and Stefan-Boltzmann constant respectively. It is assumed that the variation in temperature within the flow is such that T^4 can be expanding in Taylor series about T_∞ and neglecting higher-order terms beyond the first degree in $(T - T_\infty)$, we get:

$$T^4 \approx 4T_\infty^3 T - 3T_\infty^4 \quad (8)$$

Now, differentiating Eq. (7) w. r. to y and using Eq. (8), we get:

$$\frac{\partial q_r}{\partial y} = -\frac{16T_\infty^3 \sigma^*}{3K_s} \frac{\partial^2 T}{\partial y^2} \quad (9)$$

Using Eq. (9) in Eq. (5), we obtain:

$$u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \sigma \frac{\partial^2 T}{\partial y^2} + \frac{1}{\rho c_p} \frac{16T_\infty^3 \sigma^*}{3K_s} \frac{\partial^2 T}{\partial y^2} + \frac{Q_0}{\rho c_p} (T - T_\infty) \quad (10)$$

The following similarity transformations are introduced.

$$\begin{aligned} u &= axf'(\eta), \\ v &= -\sqrt{av}f(\eta), \theta(\eta) \\ &= \frac{T-T_\infty}{T_f-T_\infty}, \eta = y\sqrt{\frac{a}{\nu}}, \\ \psi &= x\sqrt{av}f(\eta) \end{aligned} \quad (11)$$

where, the stream function ψ defined as $u = \frac{\partial\psi}{\partial y}$, $v = -\frac{\partial\psi}{\partial x}$.

Utilizing Eq. (11), the Eqns. (4)-(5) are reduced into the following 3rd and 2nd order ordinary differential equations.

$$\begin{aligned} f''' + \beta(f''^2 - ff''''') + (1+\lambda)(ff'' - f'^2) \\ - (1+\lambda)\left(M + \frac{1}{K}\right)f' = 0 \end{aligned} \quad (12)$$

$$\left(1 + \frac{4}{3}R\right)\theta'' + Pr f\theta' + Q\theta = 0 \quad (13)$$

and Eq. (6) becomes,

$$\begin{aligned} f = S, f' = b/a = \alpha, \\ \theta' = -\gamma[1 - \theta(0)] \text{ at } \eta=0 \end{aligned} \quad (14)$$

$$\begin{aligned} f' = 0, f''(\eta) = 0, \\ \theta = 0 \text{ as } \eta \rightarrow \infty \end{aligned} \quad (15)$$

where, $\beta = \lambda_1 a$, $K = \frac{\nu}{aK^*}$, $M = \frac{\sigma B_0^2}{\rho a}$, $Pr = \frac{\nu}{\sigma}$, $R = \frac{4\sigma^* T_\infty^3}{K_1 k}$, $Q = \frac{\nu Q_0}{\rho c_p}$, $S = -\frac{v_w}{\sqrt{av}}$ ($S > 0$ for suction and $S < 0$ for injection), $\alpha = \frac{b}{a}$ and $\gamma = \frac{h}{k}\sqrt{\frac{\nu}{a}}$.

The physical quantities of interest are which is defined as $Re_x^{1/2} C_f = \frac{1}{1+\lambda}(f''(0) + \beta f'(0))$.

The surface drag coefficient (C_f) and the rate of thermal energy (Nu) is defined as:

$$C_f = \frac{\tau_w}{\frac{1}{2}\rho u_w^2}, \quad Nu_x = \frac{x(q_w + q_r)}{k(T_f - T_\infty)} \quad (16)$$

where, $\tau_w = \frac{\mu}{1+\lambda} \left\{ \left(\frac{\partial u}{\partial y}\right) + \lambda_1 \left[u \frac{\partial^2 u}{\partial x \partial y} + u \frac{\partial^2 v}{\partial x^2} + v \frac{\partial^2 u}{\partial y^2} \right] \right\}$ and $q_w = -k \left(\frac{\partial T}{\partial y}\right)_{y=0}$, q_w is heat transfer from the sheet.

Substituting the value of τ_w and q_w into Eq. (16) we get the skin friction and local Nusselt number in dimensionless form as follows:

$$\begin{aligned} c_f = 2(Re_x)^{-1/2} \alpha^{-3/2} \\ \left[\frac{1}{1+\lambda} \left\{ f''(0) + \beta(f'(0)f''(0)) \right\} \right] \end{aligned} \quad (17)$$

$$Nu_x Re_x^{-1/2} = -\left(1 + \frac{4}{3}R\right)\theta'(0) \quad (18)$$

where, $Re_x = \frac{u_w(x)}{\nu}$ is the local Reynolds number.

3. NUMERICAL PROCEDURE

The non-linear ordinary differential Eqns. (12) and (13) with the boundary conditions (14) and (15) have been solved numerically utilizing fourth-order Runge-Kutta method together with shooting technique (see Ref. [37]). This technique is explained as follows:

- I. The higher-order, non-linear differential equations are reduced into a system of first order simultaneous differential equations.
- II. The resulting IVP (Initial Value Problem) is solved by utilizing a fourth order Runge-Kutta scheme (see Refs. [38, 39]).
- III. Suitable guess values are obtained by using the Newton Raphson method.
- IV. The calculations have been performed via MATLAB bvp4c.

The iteration process is repeated until the boundary conditions are satisfied.

4. RESULTS AND DISCUSSION

The problem of MHD two-dimensional heat transfer flow of a Jeffrey liquid over a shrinking/stretching sheet is studied numerically. The Jeffrey fluid parameters λ and β are chosen in the range of $0 \leq \lambda \leq 2.0$, $0 \leq \beta \leq 0.5$. The present non-Newtonian model reduced to a viscous model in the absence of Jeffrey fluid parameters λ and β . A parametric study is conducted on Newtonian and Jeffrey fluids, in order to get the pertinent parameters on various flow fields. The velocity, skin-friction coefficient, temperature, and Nusslet number results are analysed through the plotted in Figures 2-13. Table 1, gives the comparison values of Nusslet number with those of Chen [40] and Grubka and Bobba [41] in the absence of M , Q , R , λ , β , K and γ . It is noticed that the comparison shows a concurrence with the existing outcomes and thus confirms the exactness of numerical code applied in the current work.

Table 1. Comparison results of $Nu_x Re_x^{-1/2}$ for different values of Pr (in the absence of K , γ , β , M , R , Q and λ)

Pr	Chen [40]	Grubka and Bobba [41]	Present
0.01	0.00990	0.00991	0.009000
0.72	0.46315	0.46314	0.463146
1.0	0.58199	0.58199	0.581979
10.0	2.03079	2.30800	2.308008

Table 2, depicts the effect of suction/injection parameter (S), magnetic field (M) and Deborah number (β) on the skin-friction coefficient and Nusselt number along with a variation of Pr . It is seen that the skin friction coefficient decreased with increasing suction parameter ($S > 0$) and β , while the inverse pattern is seen with $S < 0$ and M . Further, an increase in β and M increases the Nusselt Number.

Figure 2 demonstrates the influence of magnetic field parameter M on the velocity profiles against η for both Newtonian and non-Newtonian fluid cases. It is observed that the increase in M reduces the fluid velocity and boundary layer thickness. Physically, an applied magnetic field generates a

reverse force (Lorentz force) in the flow field and causes to oppose the movement of liquid. Also, it is noticed that the dimensionless fluid velocity at the sheet has a higher value for a Newtonian fluid ($\lambda=\beta=0$) than to the non-Newtonian fluid ($\lambda, \beta>0$).

Table 2. Numerical Values of $Re_x^{1/2} C_f$ and $Nu_x Re_x^{-1/2}$ for different values of S, M, β and Pr (other parameters are fixed) (i.e. $K=2.0; Q=1.0; R=0.4; \alpha=0.3; \lambda=0.2; \gamma=1.0$)

S	M	β	$Pr=1.0$		$Pr=3.0$	
			$Re_x^{1/2} C_f$	$\frac{Nu_x}{Re_x^{1/2}}$	$Re_x^{1/2} C_f$	$\frac{Nu_x}{Re_x^{1/2}}$
1.0	1.0	0.2	1.45651	0.08211	1.55417	0.09483
0.5			1.25516	0.06838	1.35234	0.08970
0.0			1.02148	0.01778	1.11647	0.05410
-0.5			10.74085	0.25238	11.2166	0.15141
-1.0			5.96083	0.16669	6.21564	0.15203
0.5	0.2		1.94889	0.06830	1.94889	0.08986
	0.5		2.11414	0.06808	2.11414	0.08980
	1.0		2.36225	0.06775	2.36225	0.08970
	1.5		2.58344	0.06747	2.58344	0.08961
	0.2	0.1	1.64132	0.06819	1.64132	0.08983
		0.2	1.57243	0.06830	1.57243	0.08986
		0.3	1.51673	0.06839	1.51673	0.08989
		0.4	1.47082	0.06845	1.47082	0.08990

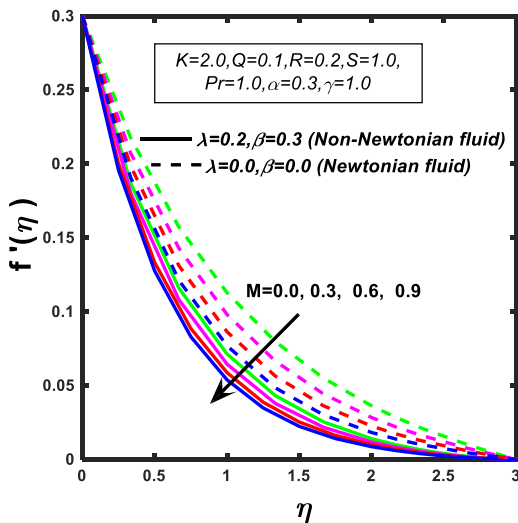


Figure 2. Effect of M on $f'(\eta)$

Figures 3(a) and 3(b) illustrate the effect of velocity ratio parameter ($\alpha=b/a$) on the velocity and temperature distributions. It is noticed that $\alpha>0$ relates to the sheet stretching while $\alpha<0$ demonstrate the shrinking sheet respectively. When $\alpha=0$ then there is no flow. It is seen that the velocity increases with the increase of α for both the stretching and shrinking cases. Also, it is observed that the rate transport of velocity curve decreases with the increasing distance η normal to the sheet and the velocity vanishes at $\eta=2$ from the sheet. Whereas α has a reverse trend in case of the temperature profile.

Figures 4(a) and 4(b) demonstrate the impact of Deborah number β on the velocity and temperature functions for both stretching and shrinking cases. The velocity becomes constant when $\alpha=0$, decreases for shrinking case and increases for stretching case. It is obvious from the figure that the momentum boundary layer converges quickly for smaller β values. Physically, small β resembles a situation where the

material has time to relax (Newtonian nature) and high β relates to non-Newtonian nature. Also, the temperature profile $\theta(\eta)$ reduced with increasing values of β for stretching case.

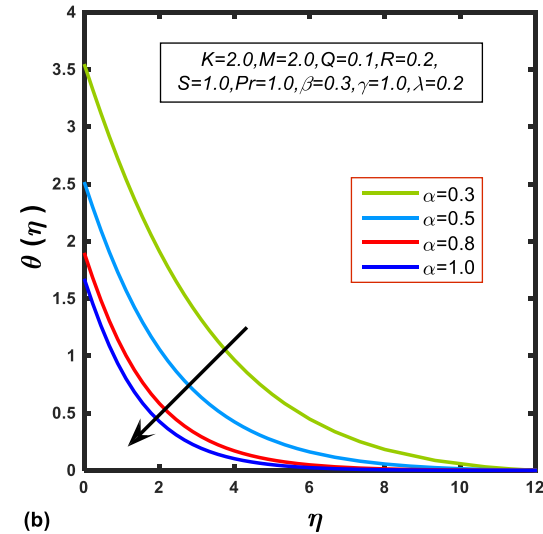
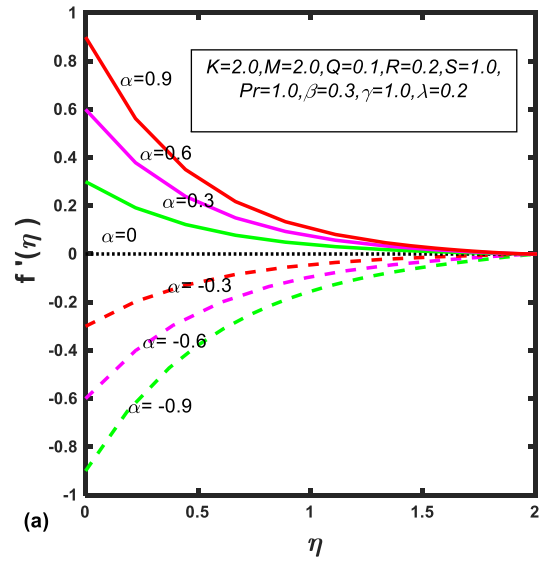
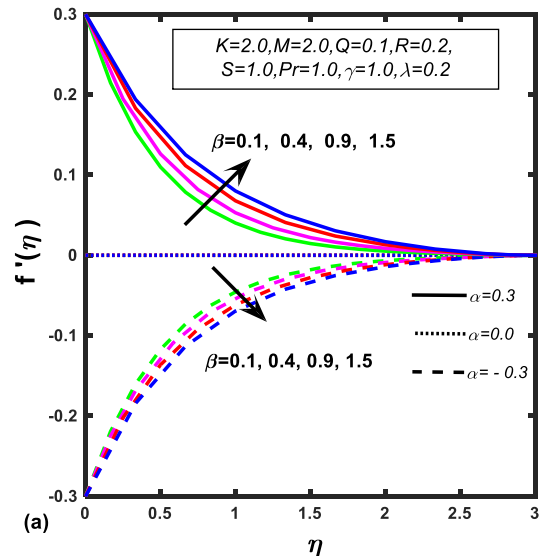


Figure 3. Effect of α on (a) $f'(\eta)$ (b) $\theta(\eta)$



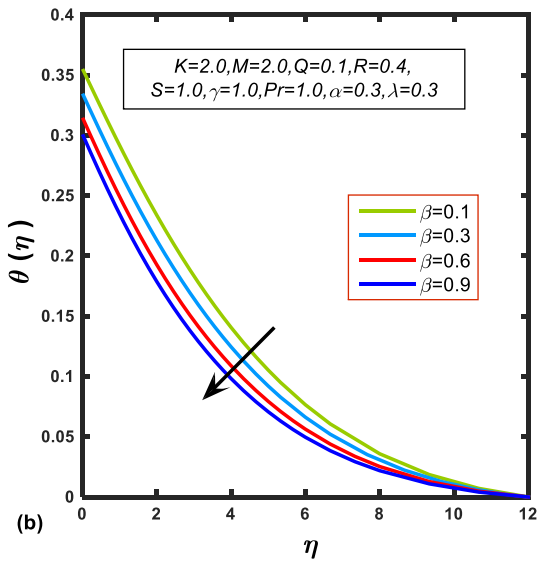


Figure 4. Effect of β on (a) $f'(\eta)$ (b) $\theta(\eta)$

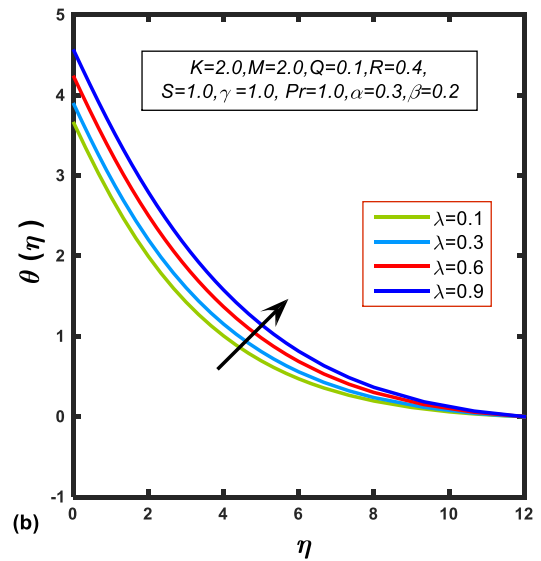


Figure 5. Effect of λ on (a) $f'(\eta)$ (b) $\theta(\eta)$

Figure 5(a) shows the influence of the ratio of relaxation and retardation times parameter λ on the velocity profiles for stretching and shrinking sheets. It is observed that the velocity decreases with the increase of λ in stretching case and reverse in shrinking case. Physically, an increase in λ means a reduction in fluid retardation time which in turn to stop the hastening of fluid motion. On the other hand, the temperature profile increases with the increase of λ which is obvious from Figure 5(b). Physically, an increase in λ leads to a rise in relaxation time and diminish in retardation time, due to this higher temperature and thicker thermal boundary layer.

The effect of permeability parameter K on the velocity profiles is shown in Figure 6. It is observed that in the case of stretching sheet ($\alpha > 0$) the velocity increases with the increase of K and β values. These results are similar to those obtained in Ref. [12]. It is also noticed from the figure that the velocity boundary layer converges quickly for small values of β .

Figures 7(a), 7(b) and 7(c) respectively show the influence of suction/injection parameter S on the velocity and temperature profiles. It is obvious from the figures that $f'(\eta)$ and $\theta(\eta)$ is found to a decrease with the increase of $S > 0$. On the other hand, inverse behaviour is seen in the injection case ($S < 0$).

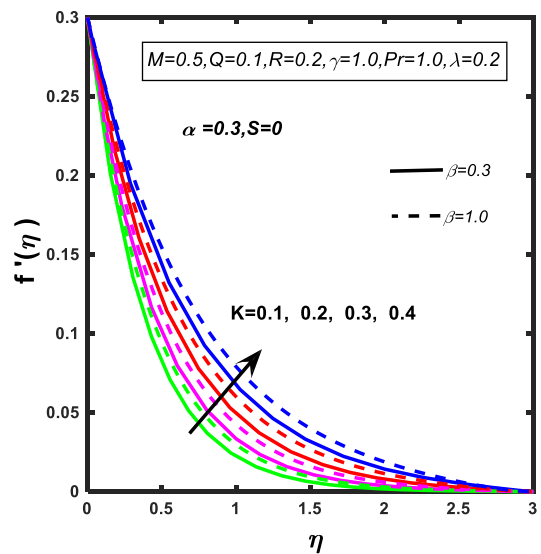
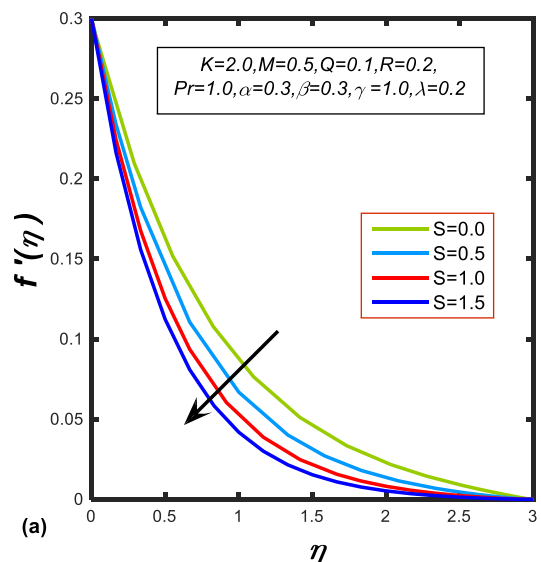
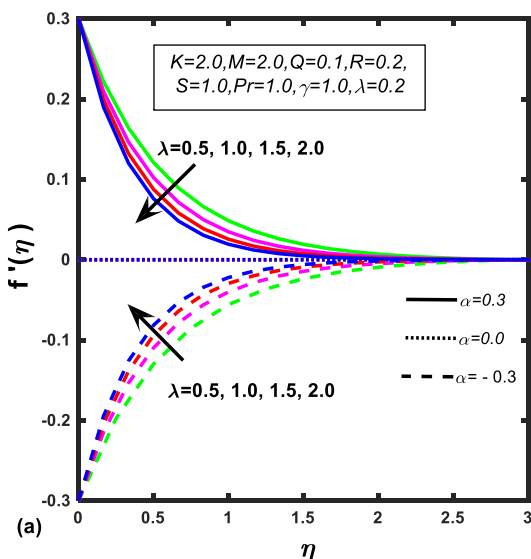


Figure 6. Effect of K on $f'(\eta)$



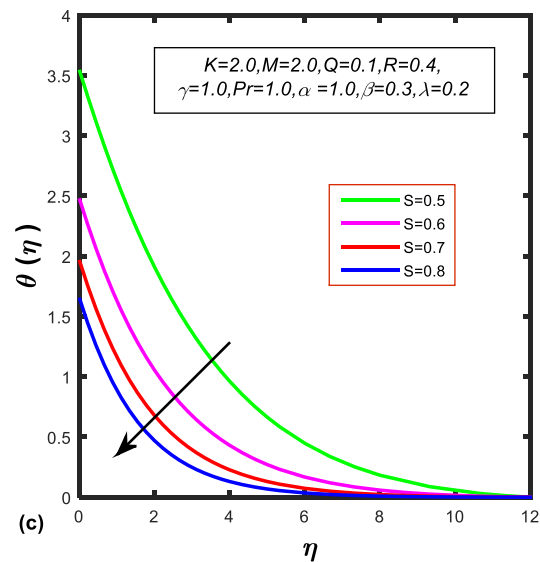
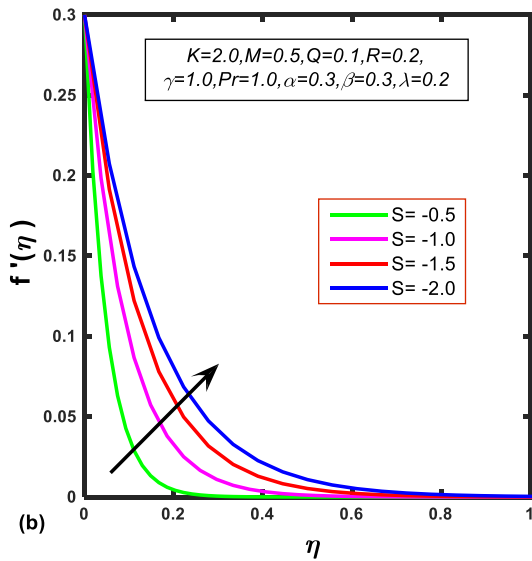


Figure 7. Effect of (a) $S \geq 0$ on $f(\eta)$ (b) $S < 0$ on $f(\eta)$ (c) $S > 0$ on $\theta(\eta)$

The influence of Biot number γ on $\theta(\eta)$ is plotted in Figure 8. It is evident that the $\theta(\eta)$ increase with the rise in γ for both stretching and shrinking cases. In general, $\gamma=0$ represents the no convective heat exchange from the surface of the sheet to the cold fluid which is far away from the sheet. For this situation, it should be understood that the surface is totally insulated. Further, it is observed that the temperature is greater in the case of shrinking when compared to the stretching sheet. These outcomes obviously reinforced from the physical perspective.

The effect of thermal radiation parameter R on the dimensionless temperature profile is offered in Figure 9. It is inspected that the fluid temperature increases by an increase in R values. This is because of the fact that as R increases; the mean absorption coefficient K_s diminishes. Hence, the divergence of radiation heat flux increases. As a result, the rate of radiative heat transfer into the liquid rises.

Figure 10 represents the temperature profile for different values of heat source/sink parameter ($Q > 0$ and $Q < 0$). It is observed that the boundary layer generates the energy, which causes the temperature profile to increase with increasing values of the heat source ($Q > 0$) whereas reverse effects are

seen in the heat sink ($Q < 0$) case. This is due to the fact that the rise of Q in the boundary layer makes energy which roots the $\theta(\eta)$ increase. From this, it is concluded that in the engineering and industrial heat transfer applications the heat sink is well appropriate for real cooling of the stretching sheet. It is also important to note that $\theta(\eta)$ decreases with increasing values of the suction parameter ($S > 0$).

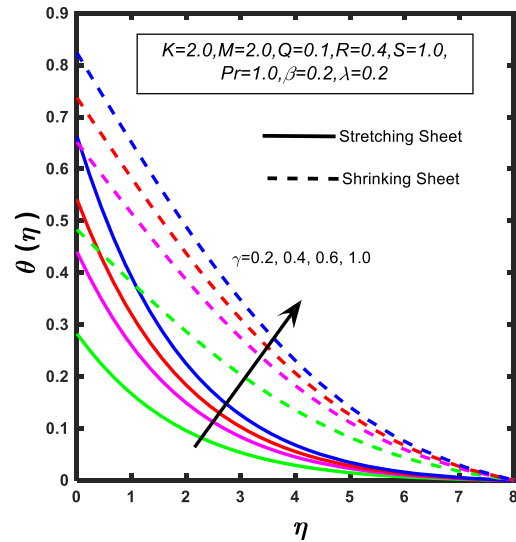


Figure 8. Effect of γ on $\theta(\eta)$

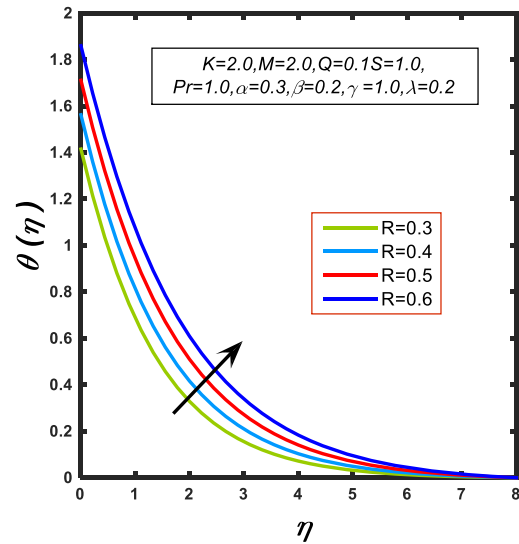


Figure 9. Effect of R on $\theta(\eta)$

Figure 11(a) portrays the attributes of skin friction coefficient concerning λ and β for stretching and shrinking cases. It is found that the skin-friction reduce in case of stretching sheet with rising β and reverse a trend in the case of a shrinking sheet. Figure 11(b) shows the influence of β on Nusselt number for stretching/shrinking sheet. The Nusselt number raised in terms of increase in β for stretching sheet, while it is reversed in case of the shrinking sheet.

Figures 12(a) and 12(b) describe the changes in Nusselt number against Prandtl number Pr for various values of Q and R respectively. It is perceived from the graphs that Nusselt number decline with the augmented values of Q and rises for Pr . The heat transfer coefficient for viscous fluid is more than that of Jeffrey fluid. In addition, the opposite trend is noticed with the rise in R .

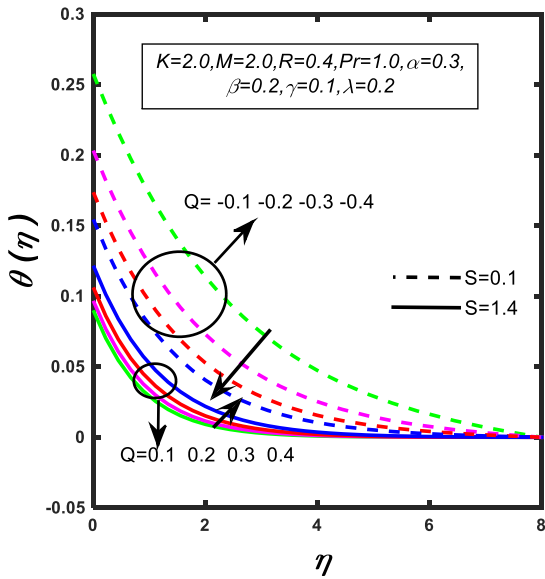
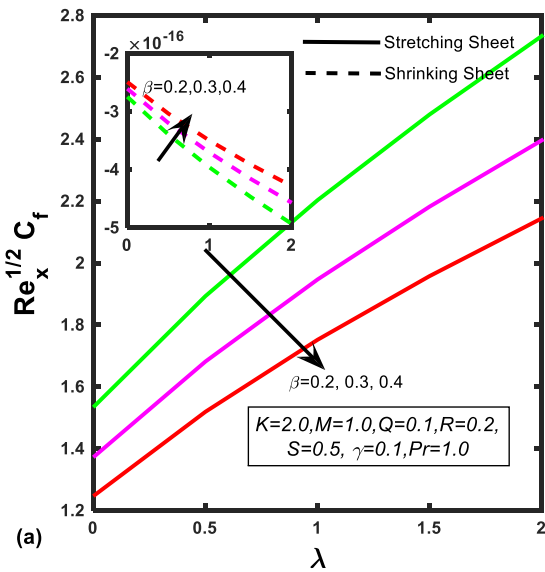
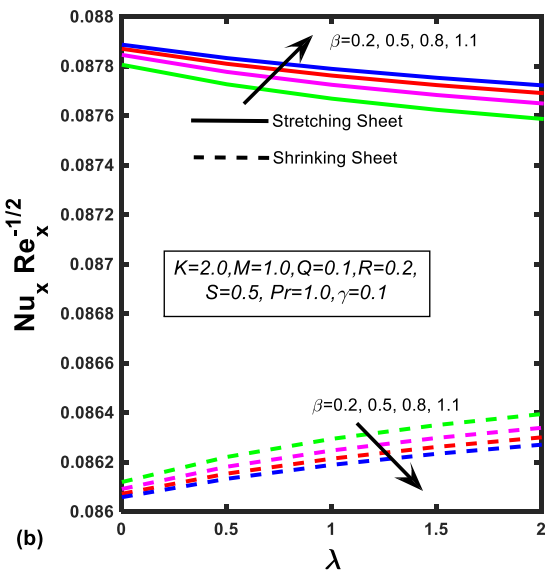


Figure 10. Effect of Q on $\theta(\eta)$

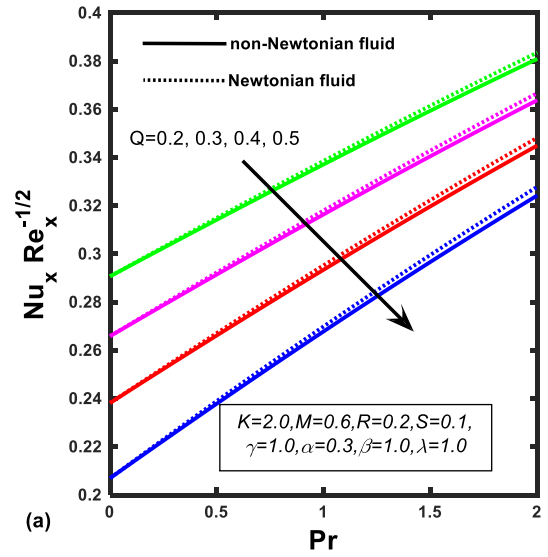


(a)

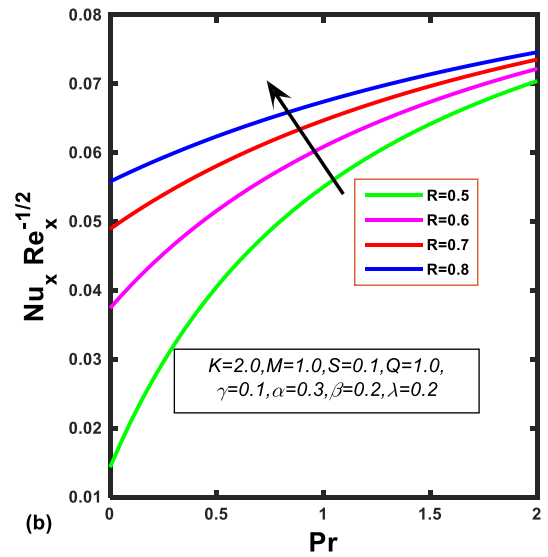


(b)

Figure 11. Effect of λ against β on (a) $Cf_x Re_x^{-1/2}$ b) $Nu_x Re_x^{-1/2}$



(a)



(b)

Figure 12. Effect of Pr against (a) Q on $Nu_x Re_x^{-1/2}$ (b) R on $Nu_x Re_x^{-1/2}$

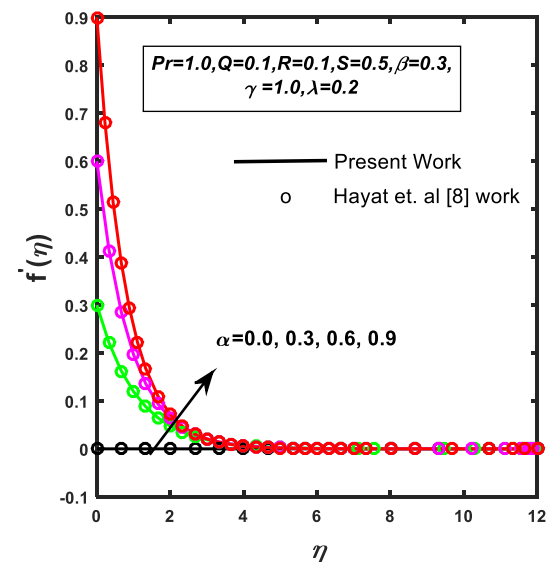


Figure 13. Comparison of the velocity distribution for various values of α between present and Hayat et al. [8] work

In order to check the accuracy of the present results, we have analyzed the outcomes of the velocity distribution for different numbers of α in Figure 13 with those reported by Hayat et al. [8] in the absence of M and K . This correlation shows an incredible arrangement.

5. CONCLUSIONS

Two-dimensional boundary layer flow of MHD heat transfer of a Jeffrey fluid with thermal radiation and heat source/sink is discussed numerically in this paper. The main points of this analysis are given below:

1. The velocity of Jeffrey fluid increases with rising stretching ratio parameter.
2. Temperature distributions have the same behaviour for different values of γ and R .
3. The variation of S on the temperature distributions is qualitatively the same when $\gamma > 0$.
4. $f'(\eta)$ enhanced for augmented values of K .
5. The velocity profile $f'(\eta)$ was declined with raising values of M for both Newtonian and non-Newtonian cases.
6. The skin friction coefficient shows reverse nature in stretching and shrinking sheet cases for diverse values of β .
7. The rate of heat transfer exhibits the same behaviour in Newtonian and non-Newtonian cases for different values of Q .

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NOMENCLATURE

a	Constant
A_1	Rivlin-Ericksen tensor
B_0	magnetic induction [T]
b	real number
c_p	specific heat at constant pressure [$J/kg/K$]
F	dimensional stream function
f	dimensionless velocity
h	convective heat transfer

K	permeability of the porous medium
k	thermal conductivity of fluid [$W/m/k$]
K_s	Rosseland mean absorption coefficient
M	magnetic field
R	radiation parameter
Nu	Nusselt number
Pr	Prandtl number
Q	heat source/sink parameter
q_r	radiative heat flux [W/m]
q_w	surface heat flux
Re_x	local Reynolds number
S	suction/injection parameter
T	fluid temperature (K)
T_f	convective fluid temperature below the moving sheet
T_w	wall temperature on a sheet at $y = 0$ (K)
T_∞	temperature far away from the wall (i)
u, v	velocity components in x-,y-directions, respectively [m/s]
u_w	Velocity of stretching sheet [ms^{-1}]
v_w	suction velocity across a stretching sheet
x	the distance along wall [m]
y	distance normal to the wall [m]

Greek symbols

α	ratio parameter
β	Deborah number
η	similarity variable
γ	Biot number
λ	ratio of relaxation and retardation times
λ_1	retardation time [s]
μ	dynamic viscosity [Pa/s]
ν	kinematic viscosity [m^2s^{-1}]
ρ	fluid density [kg/m]
σ	electric conductivity [sm^{-1}]
σ^*	Stefan-Boltzmann constant [$Wm^{-2}K^{-4}$]
θ	non-dimensional temperature
τ	Cauchy stress tensor
τ_w	shear stress along stretching sheet
ψ	stream function

Subscripts

w	sheet surface
∞	Infinity

Superscript

'	differentiation with respect to η
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