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In Depth Analysis of Stretch Function Resulting from Solving the Generalize Fractional-Order Bloch Equations Using Fractional Calculus

ABSTRACT

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In this paper the stretch function resulting from solving the fractional-order Bloch equations using fractional calculus was discussed. This function has promising results to represent diffusion signal decay from MRI images. Conventional analyses of (DWI) measurements resolve the normalized magnetization decay profiles in terms of discrete and mono-exponential components with distinct lifetimes. In complex, heterogeneous biological and biophysical samples such as tissue, multi-exponential decay functions can appear to provide truer representation to normalized magnetization decay profile than the assumption of a mono-exponential decay, but the assumption of multiple discrete components is arbitrary and is often erroneous. Moreover, interactions, between both normalized magnetization and with their environment, can result in complex normalized magnetization decay profiles that represent a continuous distribution of lifetimes. The purpose in this paper is to study different factors that influence the stretch function strength, clarity, and contrast of MRI magnetization signal relaxation by manipulating the anomalous diffusion parameters Δ , δ , G_z , β and μ . of Bloch equations. Through this study, it was found that complex normalized magnetization decay profiles behave like stretch exponential function inside power law. Further developments of this study may be useful in optimizing anomalous diffusion in tissues with neurodegenerative, and ischemic diseases.

1. INTRODUCTION

MRI and NMR application in Fractional calculus attract the attention of scientists, engineers, and researchers for a long time ago, Solving Bloch equations for various combinations of applied static, radio frequency, and gradient magnetic fields is the starting point for NMR and diffusion MRI. Significant amounts of anomalous diffusion studies have been carried out in a variety of biophysical, biological, and bioengineering complex systems for optimal sensitivity, clarity, and contrast [1-5].

An important challenge in fractional calculus science is to give a physical meaning to the fractional derivative and the resulting complex normalized magnetization decay profiles. A better way to develop this physical meaning is by studying the behavior of complex systems under known parameters like Δ , δ , G_z , β and μ . Here the dynamic anomalous diffusion models with Δ , δ , G_z , β and μ parameters resulting from solving the generalize fractional-order Bloch equations using fractional calculus become more complex as they attempt to correlate data with a multiplicity of tissue compartments, complexity, heterogeneous structure, and function [6-11]. So we expand the analysis using the Bloch equation from single exponential to multi-exponential behavior, or even to stretch exponential and power law function and from single parameter diffusion to multicompartmental diffusion and diffusion tensor imaging as well as the resulting fractional derivative related diffusion parameters [12-16].

The stretch function Figure 1 is frequently used as a purely

empirical decay law in studies of the relaxation of complex biological and biophysical systems. The stretch function can be used to describe magnetization relaxation in NMR and written as:

$$I(t) = e^{-\left(\frac{t}{\tau_0}\right)^{\beta}} \tag{1}$$

where, $0 < \beta \le 1$, and τ_0 is a parameter with the dimensions of time. This simple and relatively flexible function has been indeed successfully used in various complex biological, biophysical and biomedical fields, and it deserves thus special attention.

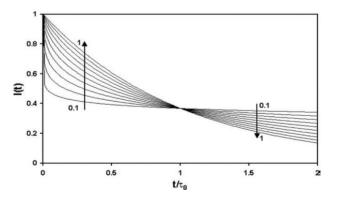


Figure 1. Plot of I (t) verses (t/τ_0) The stretched exponential decay function for several values of β (0.1(bottom curve), 0.2, ..., 0.9, 1)

The resulting solution to the Bloch-Torrey equation for the magnetization in the transverse plane derived from Magin et al. [1] can be written in the form of stretch function as:

$$M_{xy} = M_0 e^{-\left(D \mu^{2(\beta-1)} (\gamma G_z \delta)^{2\beta} \left(\Delta - \frac{2\beta - 1}{2\beta + 1} \delta\right)\right)}$$
(2)

By manipulating with the parameter, we can write this function as:

$$M_{xy} = M_0 e^{-\left(\frac{D_{\beta}}{\beta \mu_{\beta}^2} (\gamma G_z \delta \beta \mu_{\beta})^{2\beta} \left(\Delta - \frac{2\beta - 1}{2\beta + 1} \delta\right)\right)}$$
(3)

where, $0.5 < \beta < 1$, $D_{\beta} = D/\beta$, $\mu_{\beta} = \mu/\beta$, And D_{β} has unit of m^2/s , $\left(\frac{\text{meter}^2}{\text{second}}\right)$, μ_{β} has unit of *m*, (*meter*), Δ , δ has unit of *s*, (second), γ has unit of Hz/T $\left(\frac{\text{Hartes}}{\text{Tasla}}\right)$, G_z has unit of *T/meter*, $\left(\frac{\text{Tasla}}{\text{meter}}\right)$.

For **S**factor = $(\gamma G_z \delta \beta \mu_\beta)^2$ this function is unitless and we can write this function as:

$$M_{xy} = M_0 e^{-\left(\frac{D_{\beta}}{\beta\mu_{\beta}^2} (\mathbf{S}factor)^{\beta} \left(\Delta - \frac{2\beta - 1}{2\beta + 1}\delta\right)\right)}$$
(4)

For **b**'factor = $(\gamma G_z \delta \beta \mu_\beta)^{2\beta} \left(\Delta - \frac{2\beta - 1}{2\beta + 1} \delta\right)$ this function has unit of s, (second) and we can write this function as:

$$M_{xy} = M_0 e^{-\left(\frac{D_\beta}{\beta\mu_\beta^2} \times b' factor\right)}$$
(5)

For $\overrightarrow{\mathbf{b}}$ factor = $\frac{1}{\beta\mu_{\beta}^{2}} \left(\gamma G_{z}\delta\beta\mu_{\beta}\right)^{2\beta} \left(\Delta - \frac{2\beta-1}{2\beta+1}\delta\right)$ this

function has unit of s/m^2 , $\left(\frac{second}{meter^2}\right)$ and we can write this function as:

$$M_{xy} = M_0 e^{-(D_\beta \times \vec{b} factor)}$$
(6)

For **b**factor = $(\gamma G_z \delta)^2 \Delta$ and $\Delta \gg \delta$ this function has unit of s/m^2 , $\left(\frac{second}{meter^2}\right)$ and we can write this function as:

$$M_{xy} = M_0 e^{-\left(\frac{D_{\beta}}{\beta \mu_{\beta}^2} (\gamma G_z \delta \beta \mu_{\beta})^{2\beta} \Delta \left(\frac{\mathbf{b} factor}{(\gamma G_z \delta)^2 \Delta}\right)\right)}$$
(7)

$$M_{xy} = M_0 e^{-\left(\frac{D_{\beta}}{\beta\mu_{\beta}^2} (\gamma G_z \delta)^{2\beta-2} (\beta \mu_{\beta})^{2\beta} \times bfactor\right)}$$
(8)

For \mathbf{b}^* factor = $(\gamma G_z \delta)^2 \left(\Delta - \frac{2\beta - 1}{2\beta + 1} \delta\right)$ this function has unit of s/m², $\left(\frac{\text{second}}{\text{meter}^2}\right)$ and we can write this function as:

$$M_{xy} = M_0 e^{-\left(\frac{D_\beta}{\beta\mu_\beta^2}(\gamma G_z \delta)^{2\beta-2} (\beta\mu_\beta)^{2\beta} \times \mathbf{b}^* \text{ factor}\right)}$$
(9)

However, the closest form to stretched function is the first one where **S**factor = $(\gamma G_z \delta \beta \mu_\beta)^2$ this function is unitless and if **X**factor = $\frac{D_\beta}{\beta \mu_\beta^2} \left(\Delta - \frac{2\beta - 1}{2\beta + 1} \delta \right)$ this function is also unitless and we can write this function as:

$$M_{xy} = M_0 e^{-(Xfactor (Sfactor)^{\beta})}$$
(10)

Then;

$$M_{xy} = M_0 \left(e^{-(Sfactor)^{\beta}} \right)^{Xfactor}$$
(11)

2. DISCUSSION AND RESULTS

The stretched exponential function results from solving the Bloch–Torrey equation using fractional calculus provides a mechanism for introducing stretched function dynamics "tissue heterogeneity". In this section we present theoretical results for stretched-exponential function that can be applied to diffusion-weighted images. In the theoretical study, the derived magnetization attenuation curves are compared with the classical result and a more recent expression derived using stretched function models.

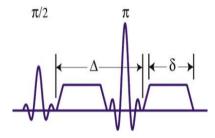


Figure 2. 2 Stejskal/Tanner diffusion preparation: Initial 90° RF pulse, followed by first diffusion lobe, 180° refocusing

RF pulse, then second diffusion lobe. Leading edges of diffusion lobes are separated by Δ . Each lobe has duration δ

The theoretical curves were plotted versus the gradient parameter, **b**factor, **b**'factor, **b** factor, **S**factor, **b*** factor and G_z for selected values of Δ , δ , G_z , β and μ Figure 2. As an example of the behavior expected, the decay of the normalized magnetization $(\frac{M_{XY}}{M_0})$, as given in Eq. (2), is plotted in Figure 3 versus b_{factor}^* where $b_{\text{factor}}^* = 0$ to $180 \ s/mm^2$ and $b_{\text{factor}}^* = (\gamma G_z \delta)^2 \left(\Delta - \frac{2\beta - 1}{2\beta + 1} \delta\right)$ for different values of D_β in the range from $D_\beta = 0.8333 \times 10^{-3}mm^2/s$ (bottom curve) to $D_\beta = 3.333 \times 10^{-3}mm^2/s$ in steps of $0.8333 \times 10^{-3}mm^2/s$ ($\Delta = 40 \times 10^{-3}s$, $\mu = 5\mu$ m, $\beta = 0.6$, $\delta = 1 \times 10^{-3}s$, $\gamma = 42.58 \ \text{MHz/T}$). We observed that as D_β increases the normalized magnetization curve change from heavy tailed decay to a straight line which strongly resembles the behavior recorded in restricted diffusion. In Figure 4 normalized magnetization M_{XY} / M_0 is plotted versus b_{factor}^* where 17

= 0 to 180 s/mm² and b_{factor}^* for different values of μ_{β} in the range from $\mu_{\beta} = 3.333 \mu m$ (bottom curve) to $\mu_{\beta} =$ 16.6667 μ m in steps of 3.333 μ m ($\Delta = 40 \times 10^{-3}$ s, D = $1 \times 10^{-3} \text{ mm}^2/\text{s}$, $G_z = 0$ to 1.5 T/m, $\delta = 1 \times 10^{-3}$ s, $\beta = 0.6$, $\gamma =$ 42.58 MHz/T). In Figure 5 normalized magnetization M_{xy} / M_0 is plotted versus b_{factor} where $b_{factor} = 0$ to 1800 s/mm^2 and $b_{factor} = (\gamma G_z \delta)^2 \Delta$ and $\Delta \gg \delta$ for different values of D_{\beta} in the range from D_{\beta} = 0.8333 $\times 10^{-3} mm^2/\text{s}$ (bottom curve) to D_{\beta} = 3.333 $\times 10^{-3} mm^2/\text{s}$ in steps of 0.8333 $\times 10^{-3} mm^2/\text{s}$ ($\Delta = 40 \times 10^{-3} s$, $\mu = 5\mu m$, $\beta = 0.6$, $δ=1 × 10^{-3}$ s, γ=42.58 MHz/T). In Figure 6 normalized magnetization M_{xy} / M_0 is plotted versus b_{factor} where $b_{factor} = 0$ to 1800 s/mm² and b_{factor} and Δ ≫ δ for different values of $μ_β$ in the range from $μ_β = 3.333 μm$ (bottom curve) to $μ_β = 16.6667 μm$ in steps of 3.333 μm ($Δ = 40 × 10^{-3}$ s, $D = 1 × 10^{-3} mm^2/s$, $G_z = 0$ to 1.5 T/m , $δ = 1 × 10^{-3}$ s, β=0.6, γ=42.58 MHz/T).

In this part we observe that as the value of μ_{β} increases the contribution of restricted diffusion increase for a fixed value of β . We can see this behavior when Eq. (2) is written either in terms of a single exponential decay, $e^{-bD_{apl}}$, where: $DapI = \frac{D}{((\gamma G_z \delta)\mu)^{2(1-\beta)}}$; or when Eq. (2) is written as a stretched exponential, $e^{-(bD_F)\beta}$, where $D_F^{\beta} = D(\Delta/\mu^2)^{1-\beta}$. Also when $\mu = \sqrt{D\Delta}$ it can be shown that the exponential form is "stretched exponential" result, $e^{-(bD)\beta}$ considered by Bennett [12, 17-20]. In Figure 5 when $\mu_{\beta} = 7.07/0.6 \ \mu m$, we observed a decrease in the apparent diffusion coefficient D_{β} as the values of β decrease and μ_{β} increase.

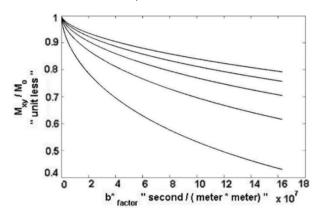


Figure 3. Stretched exponential model curves plot: M_{xy} / M_0 Versus b_{factor}^* where $b_{factor}^* = 0$ to 18×10^7 s/m² and $b_{factor}^* = (\gamma G_z \delta)^2 \left(\Delta - \frac{2\beta - 1}{2\beta + 1} \delta\right)$ for different values of D_{β} in the range from $D_{\beta} = 0.8333 \times 10^{-3} \text{mm}^2$ /s (bottom curve) to $D_{\beta} = 3.333 \times 10^{-3} \text{mm}^2$ /s in steps of $0.8333 \times 10^{-3} \text{mm}^2$ /s ($\Delta = 40 \times 10^{-3}$ s, $\mu = 5\mu$ m, $\beta = 0.6$, $\delta = 1 \times 10^{-3}$ s, $\gamma = 42.58$ MHz/T)

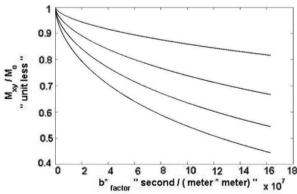


Figure 4. Stretched exponential model curves plot: M_{xy} / M_0 versus b_{factor}^* where $b_{factor}^* = 0$ to $18 \times 10^7 \text{ s/m}^2$ and $b_{factor}^* = (\gamma G_z \delta)^2 \left(\Delta - \frac{2\beta - 1}{2\beta + 1} \delta\right)$ for different values of μ_{β} in the range from $\mu_{\beta} = 3.333 \mu \text{m}$ (bottom curve) to $\mu_{\beta} = 16.6667 \mu \text{m}$ in steps of $3.333 \mu \text{m}$ ($\Delta = 40 \times 10^{-3} \text{s}$, $D = 1 \times 10^{-3} \text{mm}^2/\text{s}$, $G_z = 0$ to 1.5 T/m, $\delta = 1 \times 10^{-3} \text{s}$, $\beta = 0.6$, $\gamma = 42.58 \text{ MHz/T}$)

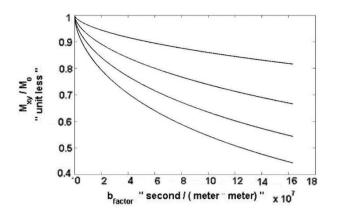


Figure 5. Stretched exponential model curves plot: M_{xy} / M_0 Versus b_{factor} where $b_{factor} = 0$ to $18 \times 10^8 \text{ s/m}^2$ and $b_{factor} = (\gamma G_z \delta)^2 \Delta$ and $\Delta \gg \delta$ for different values of D_β in the range from $D_\beta = 0.8333 \times 10^{-3} \text{mm}^2/\text{s}$ (bottom curve) to $D_\beta = 3.333 \times 10^{-3} \text{mm}^2/\text{s}$ in steps of $0.8333 \times 10^{-3} \text{mm}^2/\text{s}$ ($\Delta = 40 \times 10^{-3} \text{s}$, $\mu = 5 \mu \text{m}$, $\beta = 0.6$, $\delta = 1 \times 10^{-3} \text{s}$, $\gamma = 42.58 \text{ MHz/T}$)

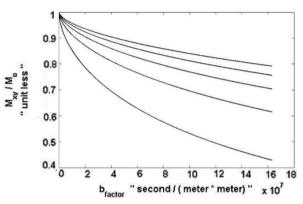


Figure 6. Stretched exponential model curves plot: M_{xy} / M_0 versus b_{factor} where $b_{factor} = 0$ to 18×10^7 s/m² and $b_{factor} = (\gamma G_z \delta)^2 \Delta$ and $\Delta \gg \delta$ for different values of μ_{β} in the range from $\mu_{\beta} = 3.333 \mu m$ (bottom curve) to $\mu_{\beta} = 16.6667 \mu m$ in steps of $3.333 \mu m$ ($\Delta = 40 \times 10^{-3}$ s, $D = 1 \times 10^{-3} mm^2/s$, $G_z = 0$ to 1.5 T/m, $\delta = 1 \times 10^{-3}$ s, $\beta = 0.6$, $\gamma = 42.58$ MHz/T)

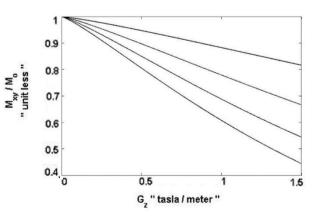


Figure 7. Stretched Exponential Model Curves Plot: M_{xy} / M_0 Versus G_z where $G_z = 0$ to 1.5 T/m for different values of D_β in the range from $D_\beta = 0.8333 \times 10^{-3} mm^2/s$ (bottom curve) to $D_\beta = 3.333 \times 10^{-3} mm^2/s$ in steps of $0.8333 \times 10^{-3} mm^2/s$ ($\Delta = 40 \times 10^{-3}s$, $\mu = 5\mu m$, $\beta = 0.6$, $\delta = 1 \times 10^{-3}s$, $\gamma = 42.58$ MHz/T)

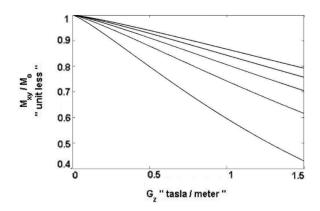


Figure 8. Stretched exponential model curves plot: M_{xy} / M_0 Versus G_z where $G_z = 0$ to 1.5 T/m for different values of μ in the range from $\mu_\beta = 3.333\mu$ m (bottom curve) to $\mu_\beta = 16.6667\mu$ m in steps of 3.333μ m ($\Delta = 40 \times 10^{-3}$ s, $D = 1 \times 10^{-3}$ mm²/s, $G_z = 0$ to 1.5 T/m, $\delta = 1 \times 10^{-3}$ s, $\beta = 0.6$, $\gamma = 42.58$ MHz/T)

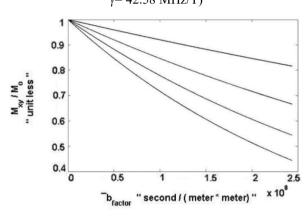


Figure 9. Stretched exponential model curves plot: M_{xy} / M_0 Versus \overline{b} factor where \overline{b} factor = 0 to 2.5 × 10⁸ s/m^2 and \overline{b} factor = $\frac{1}{\beta\mu_{\beta}^2} (\gamma G_z \delta \beta \mu_{\beta})^{2\beta} (\Delta - \frac{2\beta-1}{2\beta+1} \delta)$ for different values of D_β in the range from D_β = 0.8333× 10^{-3} mm²/s (bottom curve) to D_β = 3.333× 10^{-3} mm²/s in steps of 0.8333× 10^{-3} mm²/s ($\Delta = 40 \times 10^{-3}$ s, $\mu = 5\mu$ m, $\beta = 0.6$, $\delta = 1 \times 10^{-3}$ s, $\gamma = 42.58$ MHz/T)

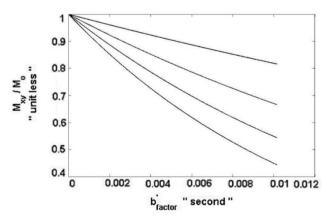


Figure 10. Stretched exponential model curves plot: M_{xy} / M_0 Versus b'_{factor} where $b'_{factor} = 0$ to 0.01 second and $b'_{factor} = (\gamma G_z \delta \beta \mu_\beta)^{2\beta} \left(\Delta - \frac{2\beta - 1}{2\beta + 1} \delta\right)$ for different values of D_β in the range from $D_\beta = 0.8333 \times 10^{-3} \text{mm}^2/\text{s}$ (bottom curve) to $D_\beta = 3.333 \times 10^{-3} \text{mm}^2/\text{s}$ in steps of 0.8333× $10^{-3} \text{mm}^2/\text{s}$ ($\Delta = 40 \times 10^{-3} \text{s}$, $\mu = 5\mu \text{m}$, $\beta = 0.6$, $\delta = 1 \times 10^{-3} \text{s}$, $\gamma = 42.58 \text{ MHz/T}$)

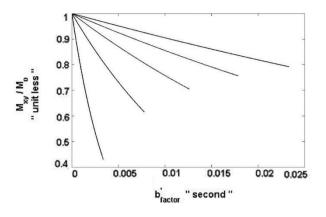


Figure 11. Stretched exponential model curves plot: M_{xy} / M_0 versus b'_{factor} where $b'_{factor} = 0$ to 0.025 second and $b'_{factor} = (\gamma G_z \delta \beta \mu_\beta)^{2\beta} \left(\Delta - \frac{2\beta - 1}{2\beta + 1} \delta\right)$ for different values of μ_β in the range from $\mu_\beta = 3.333 \mu$ m (bottom curve) to $\mu_\beta = 16.6667 \mu$ m in steps of 3.333μ m ($\Delta = 40 \times 10^{-3}$ s, $D = 1 \times 10^{-3}$ mm²/s, $G_z = 0$ to 1.5 T/m, $\delta = 1 \times 10^{-3}$ s, $\beta = 0.6$, $\gamma = 42.58$ MHz/T)

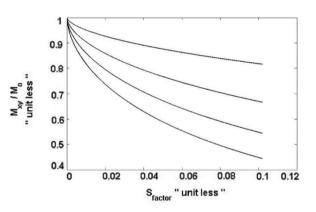


Figure 12. Stretched exponential model curves plot: M_{xy} / M_0 Versus S_{factor} where $S_{factor} = 0$ to 0.1 and $S_{factor} = (\gamma G_z \delta \beta \mu_\beta)^2$ for different values of D_β in the range from $D_\beta = 0.8333 \times 10^{-3} \text{mm}^2/\text{s}$ (bottom curve) to $D_\beta = 3.333 \times 10^{-3} \text{mm}^2/\text{s}$ in steps of $0.8333 \times 10^{-3} \text{mm}^2/\text{s}$ ($\Delta = 40 \times 10^{-3} \text{s}$, $\mu = 5 \mu \text{m}$, $\beta = 0.6$, $\delta = 1 \times 10^{-3} \text{s}$, $\gamma = 42.58$ MHz/T)

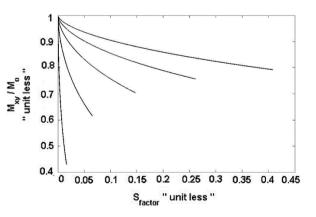


Figure 13. Stretched exponential model curves plot: M_{xy} / M_0 versus S_{factor} where $S_{factor} = 0$ to 0.45 and $S_{factor} = (\gamma G_z \delta \beta \mu_\beta)^2$ for different values of μ_β in the range from $\mu_\beta = 3.333 \mu m$ (bottom curve) to $\mu_\beta = 16.6667 \mu m$ in steps of $3.333 \mu m$ ($\Delta = 40 \times 10^{-3} s$, $D = 1 \times 10^{-3} mm^2/s$, $G_z = 0$ to 1.5 T/m , $\delta = 1 \times 10^{-3} s$, $\beta = 0.6$, $\gamma = 42.58$ MHz/T)

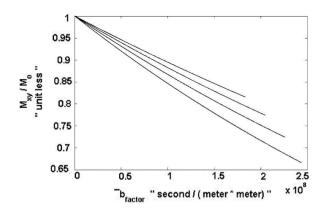


Figure 14. Stretched exponential model curves plot: M_{xy} / M_0 Versus \overline{b} factor where \overline{b} factor = 0 to 2.5 × 10⁸ s/m^2 and \overline{b} factor = $\frac{1}{\beta\mu_{\beta}^2} \left(\gamma G_z \delta\beta\mu_{\beta}\right)^{2\beta} \left(\Delta - \frac{2\beta-1}{2\beta+1} \delta\right)$ for different values of β in the range from β = 0.6 (bottom curve) to β = 0.9 in steps of 0.1 (D = 1*10⁻³ mm²/s, μ = 5 μ m, G_z = 0 to 1.5 T/m, δ = 1× 10⁻³ s, Δ =40 × 10⁻³ s, γ = 42.58 MHz/T)

In Figure 7 normalized magnetization M_{xy} / M_0 is plotted versus G_z where $G_z = 0$ to 1.5 T/m for different values of D_β in the range from $D_{\beta} = 0.8333 \times 10^{-3} mm^2/s$ (bottom curve) to $D_{\beta} = 3.333 \times 10^{-3} mm^2/s$ in steps of $0.8333 \times 10^{-3} mm^2/s$ $(\Delta = 40 \times 10^{-3} s, \mu = 5 \mu m, \beta = 0.6, \delta = 1 \times 10^{-3} s, \gamma = 42.58$ MHz/T). In Figure 8 normalized magnetization M_{xy} / M_0 is plotted versus G_z where $G_z = 0$ to 1.5 T/m for different values of μ_{β} in the range from $\mu_{\beta} = 3.333 \mu m$ (bottom curve) to $\mu_{\beta} =$ $16.6667\mu m$ in steps of $3.333\mu m$ ($\Delta = 40 \times 10^{-3} s$, D = 1×10^{-3} mm²/s, $G_z = 0$ to 1.5 T/m , $\delta = 1 \times 10^{-3}$ s, $\beta = 0.6$, $\gamma =$ 42.58 MHz/T). In Figure 9 normalized magnetization M_{xy} / M_0 is plotted versus \overline{b} factor where \overline{b} factor = 0 to 250 s/mm^2 and \overline{b} factor = $\frac{1}{\beta\mu_B^2} \left(\gamma G_z \delta\beta\mu_\beta\right)^{2\beta} \left(\Delta - \frac{2\beta-1}{2\beta+1}\delta\right)$ for different values of D_{β} in the range from $D_{\beta} = 0.8333 \times$ $10^{-3}mm^2$ /s (bottom curve) to $D_{\beta} = 3.333 \times 10^{-3}mm^2$ /s in steps of $0.8333 \times 10^{-3} mm^2/s$ ($\Delta = 40 \times 10^{-3} s$, $\mu = 5 \mu m$, $\beta =$ 0.6, $\delta = 1 \times 10^{-3}$ s, $\gamma = 42.58$ MHz/T). In Figure 10 normalized magnetization M_{xy} / M_0 is plotted versus b'_{factor} where $b'_{factor} = 0$ to 0.01 second and $b'_{factor} = (\gamma G_z \delta \beta \mu_\beta)^{2\beta} (\Delta - \beta \sigma_z \delta \beta \mu_\beta)^{2\beta}$ $\frac{2\beta-1}{2^{\beta+1}}\delta$ for different values of D_{β} in the range from $D_{\beta} =$ $0.8333 \times 10^{-3} mm^2$ /s (bottom curve) to $D_{\beta} = 3.333 \times$ $10^{-3}mm^2$ /s in steps of $0.8333 \times 10^{-3}mm^2$ /s ($\Delta = 40 \times 10^{-3}s$, $\mu = 5\mu m$, $\beta = 0.6$, $\delta = 1 \times 10^{-3} s$, $\gamma = 42.58$ MHz/T). In Figure 1 stretched exponential model curves plot, M_{xy} / M_0 versus \mathbf{b}'_{factor} where $\mathbf{b}'_{factor} = 0$ to 0.025 second and $\mathbf{b}'_{factor} = (\gamma G_z \delta \beta \mu_\beta)^{2\beta} \left(\Delta - \frac{2\beta - 1}{2\beta + 1} \delta\right)$ for different values of μ_β in the range from $\mu_{B} = 3.333 \mu m$ (bottom curve) to $\mu_{B} = 16.6667 \mu m$ in steps of $3.333 \mu m$ ($\Delta = 40 \times 10^{-3}$ s, $D = 1 \times 10^{-3} mm^2/s$, $G_z = 0$ to 1.5 T/m, $\delta = 1 \times 10^{-3}$ s, $\beta = 0.6$, $\gamma = 42.58$ MHz/T).

In Figure 13 normalized magnetization M_{xy} / M_0 is plotted versus b'_{factor} where $b'_{factor} = 0$ to 0.025 second for different values of μ_{β} in the range from $\mu_{\beta} = 3.333 \mu m$ (bottom curve) to $\mu_{\beta} = 16.6667 \mu m$ in steps of $3.333 \mu m$ ($\Delta = 40 \times 10^{-3} s$, D = $1 \times 10^{-3} mm^2/s$, $G_z = 0$ to 1.5 T/m, $\delta = 1 \times 10^{-3} s$, $\beta = 0.6$, $\gamma =$ 42.58 MHz/T). In Figure 12 normalized magnetization M_{xy} / M_0 is plotted versus S_{factor} where $S_{factor} = 0$ to 0.1 and $S_{factor} = (\gamma G_z \delta \beta \mu_\beta)^2$ for different values of D_β in the range from $D_\beta = 0.8333 \times 10^{-3} mm^2$ /s (bottom curve) to $D_\beta =$ $3.333 \times 10^{-3} mm^2$ /s in steps of $0.8333 \times 10^{-3} mm^2$ /s ($\Delta =$ $40 \times 10^{-3} s$, $\mu = 5 \mu m$, $\beta = 0.6$, $\delta = 1 \times 10^{-3} s$, $\gamma = 42.58$ MHz/T). In Figure 13 normalized magnetization M_{xy} / M_0 is plotted versus S_{factor} where $S_{factor} = 0$ to 0.45 for different values of μ_β in the range from $\mu_\beta = 3.333 \mu m$ (bottom curve) to $\mu_\beta =$ $16.6667 \mu m$ in steps of $3.333 \mu m$ ($\Delta = 40 \times 10^{-3} s$, D = $1 \times 10^{-3} mm^2/s$, $G_z = 0$ to 1.5 T/m, $\delta = 1 \times 10^{-3} s$, $\beta = 0.6$, $\gamma =$ 42.58 MHz/T). We observed that the normalized magnetization curve change from heavy tailed decay to a straight line as D_β increases and μ_β increases. Which strongly resembles the behavior of restricted diffusion.

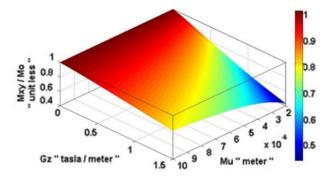


Figure 15. Stretched exponential model surface plot: M_{xy} / M_o Versus G_z where $G_z = 0$ to 1.5 T/m and μ where $\mu = 2$ to 10 μ m. (D = 1×10⁻³ mm²/s, $\beta = 0.6$, $\delta = 1×10^{-3}$ s, $\Delta = 40 \times 10^{-3}$ s, $\gamma = 42.58$ MHz/T)

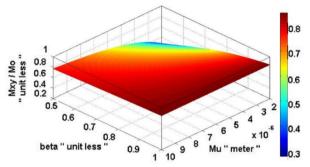


Figure 16. Stretched exponential model surface plot: M_{xy} / M_o Versus β where $\beta = 0.5$ to 1 and μ where $\mu = 2$ to 10 μ m. (D = 1×10⁻³mm²/s, G_z = 1.5 T/m, $\delta = 1 \times 10^{-3}$ s, $\Delta = 40 \times 10^{-3}$ s, $\gamma = 42.58$ MHz/T)

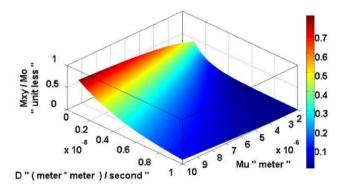


Figure 17. Stretched exponential model surface plot: M_{xy} / M_o Versus D where D = 1 × 10⁻³ to 10 × 10⁻³ mm²/s and μ where μ = 2 to 10 μ m. (β = 0.6, G_z = 1.5 T/m, δ = 1×10⁻³ s, Δ = 40×10⁻³ s, γ = 42.58 MHz/T)

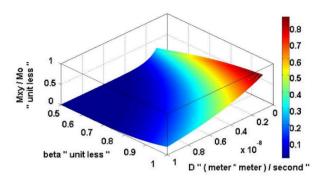


Figure 18. Stretched exponential model surface plot: M_{xy}/M_o Versus β where $\beta = 0.5$ to 1 and D where D = 1 $\times 10^{-3}$ to 10×10^{-3} mm²/s. ($\mu = 2 \times 10^{-6}$ m, $G_z = 1.5$ T/m, $\delta = 1 \times 10^{-3}$ s, $\Delta = 40 \times 10^{-3}$ s, $\gamma = 42.58$ MHz/T)

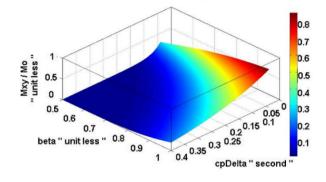


Figure 19. Stretched exponential model surface plot: M_{xy} / M_o Versus β where $\beta = 0.5$ to 1 and Δ where $\Delta = 40$ $\times 10^{-3}$ to 400×10^{-3} s. ($\mu = 2 \times 10^{-6}$ m, $G_z = 1.5$ T/m, $\delta = 1 \times 10^{-3}$ s, $D = 1 \times 10^{-3}$ mm²/s, $\gamma = 42.58$ MHz/T)

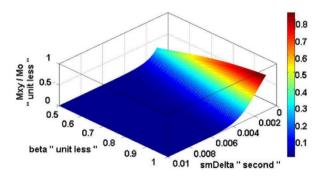


Figure 20. Stretched exponential model surface plot: M_{xy} / M_o Versus β where $\beta = 0.5$ to 1 and δ where $\delta = 1$ $\times 10^{-3}$ to 10×10^{-3} s. ($\mu = 2 \times 10^{-6}$ m, $G_z = 1.5$ T/m, $\Delta = 40 \times 10^{-3}$ s, $D = 1 \times 10^{-3}$ mm²/s, $\gamma = 42.58$ MHz/T)

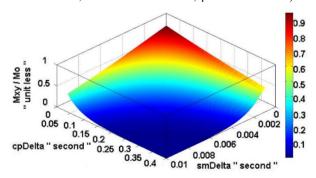


Figure 21. Stretched exponential model surface plot: M_{xy} / M_o Versus Δ where $\Delta = 40 \times 10^{-3}$ to 400×10^{-3} s and δ where $\delta = 1 \times 10^{-3}$ to 10×10^{-3} s. ($\mu = 2 \times 10^{-6}$ m, $G_z = 1.5$ T/m, $\beta = 0.6$, $D = 1 \times 10^{-3}$ mm²/s, $\gamma = 42.58$ MHz/T)

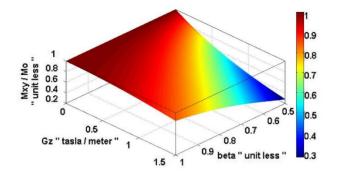


Figure 22. Stretched exponential model surface plot M_{xy} / M_o Versus G_z where $G_z = 0$ to 1.5 T/m and β where $\beta = 0.5$ to 1. ($\mu = 2 \times 10^{-6}$ m, $\Delta = 40 \times 10^{-3}$ s, $\delta = 1 \times 10^{-3}$ s, $D = 1 \times 10^{-3}$ mm²/s, $\gamma = 42.58$ MHz/T)

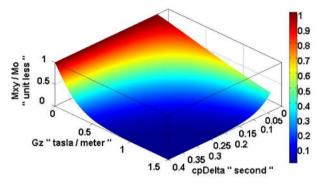


Figure 23. Stretched exponential model surface plot: M_{xy} / M_o Versus G_z where $G_z = 0$ to 1.5 T/m and Δ where Δ $= 40 \times 10^{-3}$ to 40×10^{-3} s. ($\mu = 2 \times 10^{-6}$ m, $\beta = 0.6$, $\delta =$ 1×10^{-3} s, $D = 1 \times 10^{-3}$ mm²/s, $\gamma = 42.58$ MHz/T)

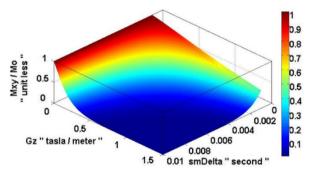


Figure 24. Stretched exponential model surface plot: M_{xy} / M_o Versus G_z where $G_z = 0$ to 1.5 T/m and δ where $\delta = 1 \times 10^{-3}$ to 10×10^{-3} s . ($\mu = 2 \times 10^{-6}$ m, $\beta = 0.6$, $\Delta = 40 \times 10^{-3}$ s, $D = 1 \times 10^{-3}$ mm²/s, $\gamma = 42.58$ MHz/T)

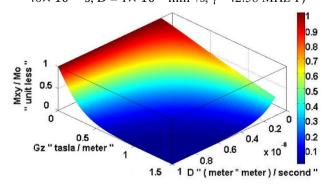


Figure 25. Stretched exponential model surface plot: M_{xy} / M_o Versus G_z where $G_z = 0$ to 1.5 T/m and D where D $= 1 \times 10^{-3}$ to 10×10^{-3} mm²/s. ($\mu = 2 \times 10^{-6}$ m, $\beta = 0.6$, $\Delta = 40 \times 10^{-3}$ s, $\delta = 1 \times 10^{-3}$ s, $\gamma = 42.58$ MHz/T)

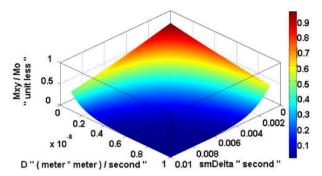


Figure 26. Stretched exponential model surface plot: M_{xy}/M_o Versus δ where $\delta = 1 \times 10^{-3}$ to 10×10^{-3} s and D where D = 1×10^{-3} to 10×10^{-3} mm²/s. ($\mu = 2 \times 10^{-6}$ m, $\beta = 0.6$, $\Delta = 40 \times 10^{-3}$ s, $G_z = 1.5$ T/m, $\gamma = 42.58$ MHz/T)

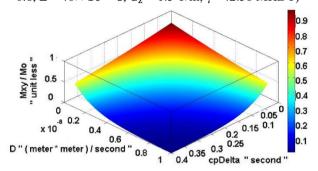


Figure 27. Stretched exponential model surface plot: M_{xy} / M_o Versus δ where $\Delta = 40 \times 10^{-3}$ to 400×10^{-3} s and D where D = 1×10^{-3} to 10×10^{-3} mm²/s. ($\mu = 2 \times 10^{-6}$ m, $\beta = 0.6$, $\delta = 1 \times 10^{-3}$ s, $G_z = 1.5$ T/m, $\gamma = 42.58$ MHz/T)

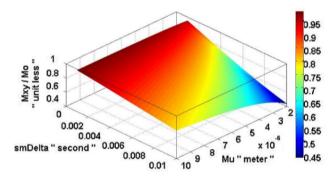


Figure 28. Stretched exponential model surface plot: M_{xy} / M_o Versus δ where $\delta = 1 \times 10^{-3}$ to 10×10^{-3} s and μ where $\mu = 2$ to $10 \ \mu$ m. (D = $1 \times 10^{-3} \text{ mm}^2/\text{s}$, $\beta = 0.6$, $\Delta = 40 \times 10^{-3}$ s, $G_z = 1.5$ T/m, $\gamma = 42.58$ MHz/T)

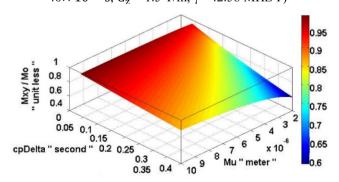


Figure 29. Stretched exponential model surface plot: M_{xy} / M_o Versus Δ where $\Delta = 40 \times 10^{-3}$ to 400×10^{-3} s and μ where $\mu = 2$ to 10 μ m. (D = 1×10⁻³ mm²/s, $\beta = 0.6$, $\delta = 1 \times 10^{-3}$ s, $G_z = 1.5$ T/m, $\gamma = 42.58$ MHz/T)

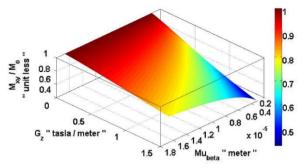


Figure 30. Stretched exponential model surface plot: M_{xy}/M_o Versus G_z where $G_z = 0$ to 1.5 T/m and μ_β where $\mu_\beta = 0.333 \times 10^{-5}$ to 1.67×10^{-5} m. (D = 1× 10⁻³ mm²/s, β = 0.6, $\delta = 1 \times 10^{-3}$ s, $\Delta = 40 \times 10^{-3}$ s, $\gamma = 42.58$ MHz/T)

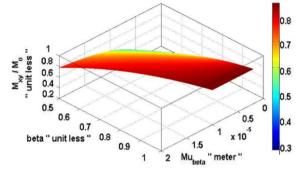


Figure 31. Stretched exponential model surface plot: M_{xy} / M_o Versus β where $\beta = 0.5$ to 1 and μ_β where μ_β =0.333× 10⁻⁵ to 1.67 × 10⁻⁵m. (D = 1× 10⁻³mm²/s, G_z = 1.5 T/m, $\delta = 1 \times 10^{-3}$ s, $\Delta = 40 \times 10^{-3}$ s, $\gamma = 42.58$ MHz/T)

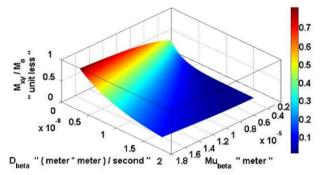


Figure 32. Stretched exponential model surface plot: M_{xy} / M_o Versus D_β where D_β = 1.7 × 10⁻³ to 16.7× 10⁻³ mm²/s, and μ_β where μ_β =3.33× 10⁻⁵ to 1.67 × 10⁻⁵m. (β = 0.6, G_z = 1.5 T/m, δ = 1× 10⁻³ s, Δ = 40× 10⁻³ s, γ= 42.58 MHz/T)

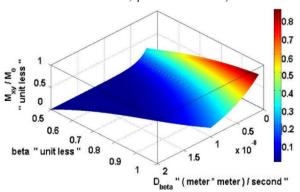


Figure 33. Stretched exponential model surface plot: M_{xy} / M_o Versus D_β where $D_\beta = 1.7 \times 10^{-3}$ to 16.7×10^{-3} mm²/s, and β where $\beta = 0.5$ to 1. ($\mu = 2\mu$ m, $G_z = 1.5$ T/m, $\delta = 1 \times 10^{-3}$ s, $\Delta = 40 \times 10^{-3}$ s, $\gamma = 42.58$ MHz/T)

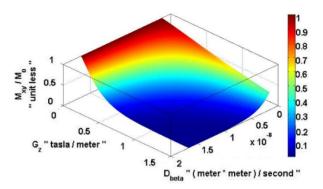


Figure 34. Stretched exponential model surface plot: M_{xy}/M_o Versus D_β where $D_\beta = 1.7 \times 10^{-3}$ to 16.7×10^{-3} mm²/s, and G_z where $G_z = 0$ to 1.5 T/m. ($\mu = 2\mu$ m, $\beta = 0.6$, $\delta = 1 \times 10^{-3}$ s, $\Delta = 40 \times 10^{-3}$ s, $\gamma = 42.58$ MHz/T)

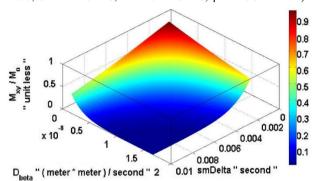


Figure 35. Stretched exponential model surface plot: M_{xy} / M_o Versus D_β where $D_\beta = 1.7 \times 10^{-3}$ to 16.7×10^{-3} mm²/s, and δ where $\delta = 1 \times 10^{-3}$ to 10×10^{-3} s. ($\mu = 2\mu$ m, $\beta = 0.6$, $G_z = 1.5$ T/m, $\Delta = 40 \times 10^{-3}$ s, $\gamma = 42.58$

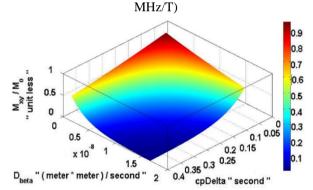


Figure 36. Stretched exponential model surface plot: M_{xy} / M_o Versus D_β where $D_\beta = 1.7 \times 10^{-3}$ to 16.7×10^{-3} mm²/s, and Δ where $\Delta = 40 \times 10^{-3}$ to 400×10^{-3} s. ($\mu = 2\mu$ m, $\beta = 0.6$, $G_z = 1.5$ T/m, $\delta = 1 \times 10^{-3}$ s, $\gamma = 42.58$ MHz/T)

Another example of the behavior expected, the decay of the normalized magnetization (M_{xy}/M_0) , is plotted in Figure 14 versus \overline{b} factor where \overline{b} factor = 0 to 250 s/mm^2 and \overline{b} factor = $\frac{1}{\beta\mu_{\beta}^2} \left(\gamma G_z \delta\beta\mu_{\beta}\right)^{2\beta} \left(\Delta - \frac{2\beta-1}{2\beta+1}\delta\right)$ for different values of β in the range from $\beta = 0.6$ (bottom curve) to $\beta = 0.9$ in steps of 0.1 (D = 1*10⁻³mm²/s, $\mu = 5\mu$ m, $G_z = 0$ to 1.5 T/m, $\delta = 1 \times 10^{-3}$ s, $\Delta = 40 \times 10^{-3}$ s, $\gamma = 42.58$ MHz/T).

In Figure 15, Eq. (2) normalized magnetization M_{xy}/M_o versus G_z was plotted where $G_z = 0$ to 1.5 T/m and μ where μ = 2 to 10 μ m. (D = 1×10⁻³mm²/s, β = 0.6, δ = 1×10⁻³s, Δ = 40 × 10⁻³ s, γ = 42.58 MHz/T). In Figure 16, Eq. (2) normalized magnetization M_{xy}/M_o versus β was plotted

where, $\beta = 0.5$ to 1 and μ where $\mu = 2$ to 10 μ m. (D = 1×10^{-3} mm²/s, G_z = 1.5 T/m, $\delta = 1 \times 10^{-3}$ s, $\Delta = 40 \times 10^{-3}$ s, γ = 42.58 MHz/T). In Figure 17, Eq. (2) normalized magnetization M_{xy}/M_o Versus D was plotted where D = 1 $\times 10^{-3}$ to 10×10^{-3} mm²/s and μ where $\mu = 2$ to $10 \ \mu$ m. ($\beta =$ 0.6, $G_z = 1.5$ T/m, $\delta = 1 \times 10^{-3}$ s, $\Delta = 40 \times 10^{-3}$ s, $\gamma = 42.58$ MHz/T). In Figure 18, Eq. (2) normalized magnetization M_{xy}/M_0 versus β was plotted where $\beta = 0.5$ to 1 and D where $D = 1 \times 10^{-3}$ to 10×10^{-3} mm²/s. ($\mu = 2 \times 10^{-6}$ m, $G_z = 1.5$ T/m, $\delta = 1 \times 10^{-3}$ s, $\Delta = 40 \times 10^{-3}$ s, $\gamma = 42.58$ MHz/T). In Figure 19, Eq. (2) normalized magnetization M_{xy}/M_o versus β was plotted where $\beta = 0.5$ to 1 and Δ where $\Delta = 40 \times 10^{-3}$ to 400×10^{-3} s. ($\mu = 2 \times 10^{-6}$ m, $G_z = 1.5$ T/m, $\delta = 1 \times 10^{-3}$ s, $D = 1 \times 10^{-3} \text{mm}^2/\text{s}$, $\gamma = 42.58 \text{ MHz/T}$). In Figure 20, Eq. (2) normalized magnetization M_{xy}/M_o versus β was plotted where $\beta = 0.5$ to 1 and δ where $\delta = 1 \times 10^{-3}$ to 10×10^{-3} s. $(\mu = 2 \times 10^{-6} \text{m}, \text{G}_z = 1.5 \text{ T/m}, \Delta = 40 \times 10^{-3} \text{s}, D = 1 \times 10^{-3} \text{s}$ 10^{-3} mm² /s, γ = 42.58 MHz/T). In Figure 21, Eq. (2) normalized magnetization M_{xy}/M_o versus Δ was plotted where $\Delta = 40 \times 10^{-3}$ to 400×10^{-3} s and δ where $\delta = 1$ $\times 10^{-3}$ to 10×10^{-3} s. ($\mu = 2 \times 10^{-6}$ m, $G_z = 1.5$ T/m, $\beta = 0.6$, $D = 1 \times 10^{-3} \text{mm}^2/\text{s}$, $\gamma = 42.58 \text{ MHz/T}$). In Figure 22 Eq. (2) normalized magnetization Mxy /Mo versus Gz was plotted where $G_z = 0$ to 1.5 T/m and β where $\beta = 0.5$ to 1. ($\mu =$ 2×10^{-6} m, $\Delta = 40 \times 10^{-3}$ s, $\delta = 1 \times 10^{-3}$ s, $D = 10^{-3}$ s, $D = 10^{-3}$ s, $D = 10^{-3}$ s, $D = 10^{-3}$ s, D = 1 10^{-3} mm² /s, γ = 42.58 MHz/T). In Figure 23 Eq. (2) normalized magnetization M_{xy}/M_o versus G_z was plotted where $G_z = 0$ to 1.5 T/m and Δ where $\Delta = 40 \times 10^{-3}$ to 40×10^{-3} s. ($\mu = 2 \times 10^{-6}$ m, $\beta = 0.6$, $\delta = 1 \times 10^{-3}$ s, $D = 1 \times 10^{-3}$ s. 10^{-3} mm² /s, γ = 42.58 MHz/T). In Figure 24 Eq. (2) normalized magnetization M_{xv}/M_o versus G_z was plotted where $G_z = 0$ to 1.5 T/m and δ where $\delta = 1 \times 10^{-3}$ to 10×10^{-3} s. ($\mu = 2 \times 10^{-6}$ m, $\beta = 0.6$, $\Delta = 40 \times 10^{-3}$ s, $D = 1 \times$ 10^{-3} mm² /s, $\gamma = 42.58$ MHz/T). In Figure 25 Eq. (2) normalized magnetization M_{xy}/M_o versus G_z was plotted where $G_z = 0$ to 1.5 T/m and D where $D = 1 \times 10^{-3}$ to 10×10^{-3} mm²/s. ($\mu = 2 \times 10^{-6}$ m, $\beta = 0.6$, $\Delta = 40 \times 10^{-3}$ s, δ = 1×10^{-3} s, $\gamma = 42.58$ MHz/T). In Figure 26 Eq. (2) normalized magnetization M_{xy}/M_o versus δ was plotted where $\delta = 1 \times 10^{-3}$ to 10×10^{-3} s and D where D = 1×10^{-3} to 10×10^{-3} mm²/s. ($\mu = 2 \times 10^{-6}$ m, $\beta = 0.6$, $\Delta =$ 40×10^{-3} s, $G_z = 1.5$ T/m, $\gamma = 42.58$ MHz/T). In Figure 27 Eq. (2) normalized magnetization M_{xy}/M_o versus δ was plotted where $\Delta = 40 \times 10^{-3}$ to 400×10^{-3} s and D where D = 1×10^{-3} to 10×10^{-3} mm²/s. ($\mu = 2 \times 10^{-6}$ m, $\beta = 0.6$, $\delta =$ 1×10^{-3} s, G_z = 1.5 T/m, γ = 42.58 MHz/T). In Figure 28 Eq. (2) normalized magnetization M_{xy}/M_o versus δ was plotted where $\delta = 1 \times 10^{-3}$ to 10×10^{-3} s and μ where $\mu = 2$ to $10 \ \mu$ m. magnetization M_{xy}/M_o versus Δ was plotted where $\Delta = 40 \times$ 10^{-3} to 400×10^{-3} s and μ where $\mu = 2$ to 10 μ m. (D = 1×10^{-3} mm²/s, $\beta = 0.6$, $\delta = 1 \times 10^{-3}$ s, $G_z = 1.5$ T/m, $\gamma =$ 42.58 MHz/T). In Figure 30 Eq. (2) normalized magnetization M_{xy}/M_0 versus G_z was plotted where $G_z = 0$ to 1.5 T/m and μ_{β} where $\mu_{\beta} = 0.333 \times 10^{-5}$ to 1.67×10^{-5} m. (D = 1 × $10^{-3}mm^2$ /s, β = 0.6, δ = 1 × 10⁻³ s, Δ = 40 × 10⁻³ s, γ = 42.58 MHz/T). In Figure 31 Eq. (2) normalized magnetization M_{xy}/M_o versus β was plotted where $\beta = 0.5$ to 1 and μ_{β} where μ_{β} =0.333 $\times\,10^{-5}$ to 1.67 $\,\times\,10^{-5}$ m. (D = 1 \times 10^{-3} mm²/s, G_z = 1.5 T/m, δ = 1×10⁻³ s, Δ = 40×10⁻³ s, γ = 42.58 MHz/T). In Figure 32 Eq. (2) normalized magnetization M_{xy}/M_o versus D_β was plotted where $D_\beta =$ 1.7×10^{-3} to 16.7×10^{-3} mm²/s, and μ_{β} where $\mu_{\beta} = 3.33 \times$ 10^{-5} to 1.67×10^{-5} m. ($\beta = 0.6$, $G_z = 1.5$ T/m, $\delta = 1 \times 10^{-3}$ s, $\Delta = 40 \times 10^{-3}$ s, $\gamma = 42.58$ MHz/T). In Figure 33 Eq. (2) normalized magnetization M_{xy}/M_o versus D_β was plotted where $D_{\beta} = 1.7 \times 10^{-3}$ to 16.7×10^{-3} mm²/s, and β where β =0.5 to 1. ($\mu = 2\mu m$, G_z = 1.5 T/m, $\delta = 1 \times 10^{-3}$ s, $\Delta =$ 40×10^{-3} s, $\gamma = 42.58$ MHz/T). In Figure 34 Eq. (2) normalized magnetization M_{xy}/M_o versus D_β was plotted where $D_{\beta} = 1.7 \times 10^{-3}$ to 16.7×10^{-3} mm²/s, and G_z where $G_z = 0$ to 1.5 T/m. ($\mu = 2\mu m$, $\beta = 0.6$, $\delta = 1 \times 10^{-3}$ s, $\Delta = 40 \times 10^{-3}$ s, $\gamma = 42.58$ MHz/T). In Figure 35 Eq. (2) normalized magnetization $M_{xy}\,/M_o$ versus D_β was plotted where $D_{\beta} = 1.7 \times 10^{-3}$ to 16.7×10^{-3} mm²/s, and δ where δ =1×10⁻³ to 10×10⁻³ s. (μ = 2 μ m, β = 0.6, G_z = 1.5 T/m, Δ = 40×10^{-3} s, γ = 42.58 MHz/T). In Figure 36 Eq. (2) normalized magnetization M_{xy}/M_o versus D_β was plotted where $D_{\beta} = 1.7 \times 10^{-3}$ to 16.7×10^{-3} mm²/s, and Δ where $\Delta = 40 \times 10^{-3}$ to 400×10^{-3} s. ($\mu = 2\mu m$, $\beta = 0.6$, $G_z = 1.5$ T/m, δ = 1 \times 10⁻³ s, γ = 42.58 MHz/T). In Figure 37 Eq. (2) normalized magnetization M_{xy}/M_o versus δ was plotted where $\delta = 1 \times 10^{-3}$ to 10×10^{-3} s and μ_{β} where $\mu_{\beta} = 0.333 \times 10^{-3}$ 10^{-5} to 1.67×10^{-5} m. (D = 1×10^{-3} mm²/s, G_z = 1.5 T/m, $\beta = 0.6$, $\Delta = 40 \times 10^{-3}$ s, $\gamma = 42.58$ MHz/T). Finally, in Figure 38 Eq. (2) the normalized magnetization M_{xy}/M_o versus Δ was plotted where $\Delta\!\!=40\!\times$ 10^{-3} to $400\!\times$ 10^{-3} s and μ_{β} where $\mu_{B} = 0.333 \times 10^{-5}$ to 1.67×10^{-5} m. (D = 1×10⁻³mm²/s, G_z = 1.5 T/m, β = 0.6, δ = 1× 10⁻³ s, γ = 42.58 MHz/T).

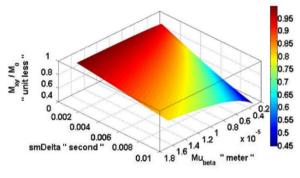


Figure 37. Stretched exponential model surface plot: M_{xy} / M_o Versus δ where $\delta = 1 \times 10^{-3}$ to 10×10^{-3} s and μ_β where $\mu_\beta = 0.333 \times 10^{-5}$ to 1.67×10^{-5} m. (D = 1× 10^{-3} mm²/s, $G_z = 1.5$ T/m, $\beta = 0.6$, $\Delta = 40 \times 10^{-3}$ s, $\gamma = 42.58$

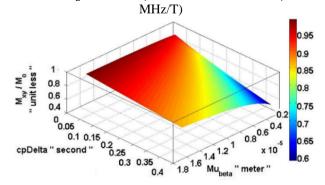


Figure 38. Stretched exponential model surface plot: M_{xy} / M_o Versus Δ where $\Delta = 40 \times 10^{-3}$ to 400×10^{-3} s and μ_{β} where $\mu_{\beta} = 0.333 \times 10^{-5}$ to 1.67×10^{-5} m. (D = 1×10^{-3} mm²/s, $G_z = 1.5$ T/m, $\beta = 0.6$, $\delta = 1 \times 10^{-3}$ s, $\gamma = 42.58$ MHz/T)

3. CONCLUSION

In this paper we discuss the stretch function resulting from solving the generalize fractional-order Bloch equations using fractional calculus in details. The theoretical curves were plotted versus gradient parameter the bfactor, b'factor, \overline{b} factor, Sfactor, b* factor and G_7 for selected values of $\Delta,\,\delta,\,G_z$, β and $\mu.$ Stretched exponential function surface plot were also plotted, different shapes of decays were observed versus different values of , bfactor, b'factor, \overline{b} factor, Sfactor, b* factor and G_{z} for selected values of Δ , δ , G_z , β and μ .

We observe that the stretch function resulting from solving the generalize fractional-order Bloch equations behaves like pure exponential function with rats D_{β} when we use \overline{b} factor, and behaves like stretch exponential function inside power law when we use S factor and X factor. Further developments of this study may be useful in optimizing anomalous diffusion in tissues with neurodegenerative, and ischemic diseases.

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